

The ring of Brownian motion: the good, the bad, and the simply silly

Peter Hänggi

Institute of Physics, University of Augsburg,
D-86135 Augsburg



Brownian motion

Gypsum cristals in a closterium moniliferum

Movie

Why you should **not** do Brownian motion

- You know nothing about the subject
- Many very good people worked on it
(Einstein, Langevin, Smoluchowski, Ornstein, Uhlenbeck, Wiener, Onsager, Stratonovich, ...)
- You don't have your own pet theory yet

Why you should do Brownian motion

- You know nothing about the subject
- Many very good people worked on it
- You still can do your own pet theory

Robert Brown (1773-1858)



Source: www.anbg.gov.au



Source: permission kindly granted by Prof. Brian J. Ford
<http://www.brianjford.com/wbbrown.htm>

1827 – irregular motion of granules of pollen in liquids

- Brown, Phil. Mag. **4**, 161 (1928)
- Deutsch: *Did Robert Brown observe Brownian Motion: probably not*, Sci. Am. **256**, 20 (1991)
- Ford: “*Brownian movement in clarkia pollen: a reprise of the first observations*”,
The Microscope **39**, 161 (1991)

Jan Ingen-Housz (1730-1799)



Source: www.americanchemistry.com

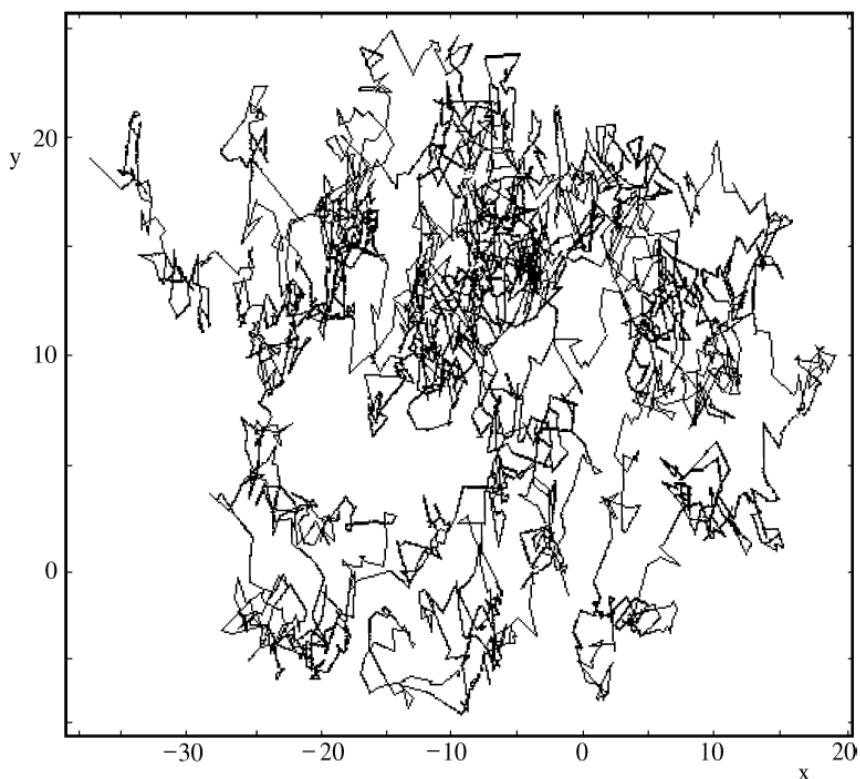


To see clearly how one can deceive one's mind on this point if one is not careful, one has only to place a drop of alcohol in the focal point of a microscope and introduce a little finely ground charcoal therein, and one will see these corpuscles in a confused, continuous and violent motion, as if they were animalcules which move rapidly around.

Mean squared displacement

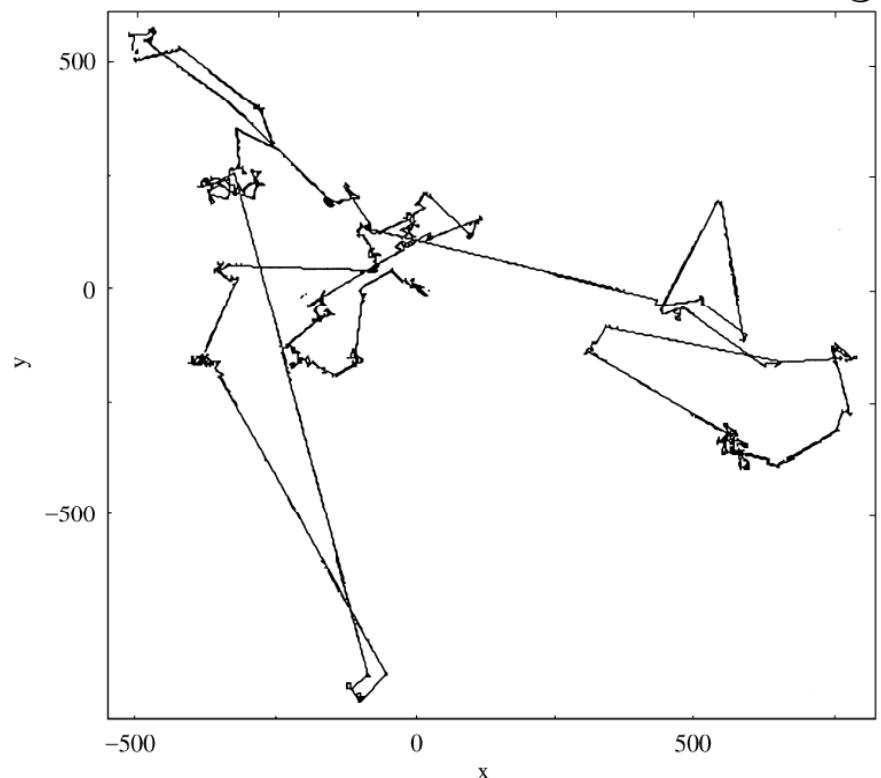
$$\langle x^2(t) \rangle \propto t^\alpha$$

Brownian movement $\alpha = 1$



Source: Physica A **282**, 13 (2000)

Lévy-Brownian movement $\alpha = \frac{4}{3}$



Source: Physica A **282**, 13 (2000)

Theory of Brownian motion

W. Sutherland (1858-1911)

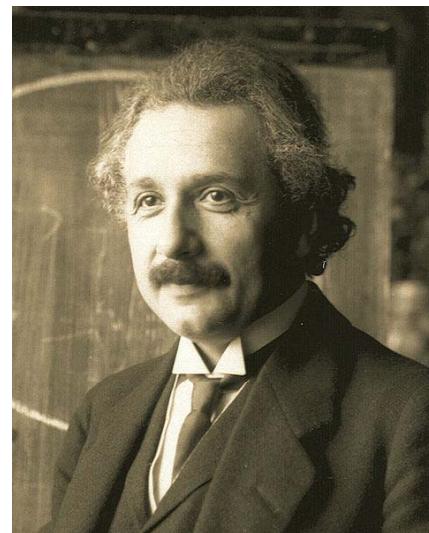


Source: www.theage.com.au

$$D = \frac{RT}{6\pi\eta aC}$$

Phil. Mag. **9**, 781 (1905)

A. Einstein (1879-1955)



Source: [wikipedia.org](https://en.wikipedia.org)

$$\langle x^2(t) \rangle = 2Dt$$

$$D = \frac{RT}{N} \frac{1}{6\pi k P}$$

Ann. Phys. **17**, 549 (1905)

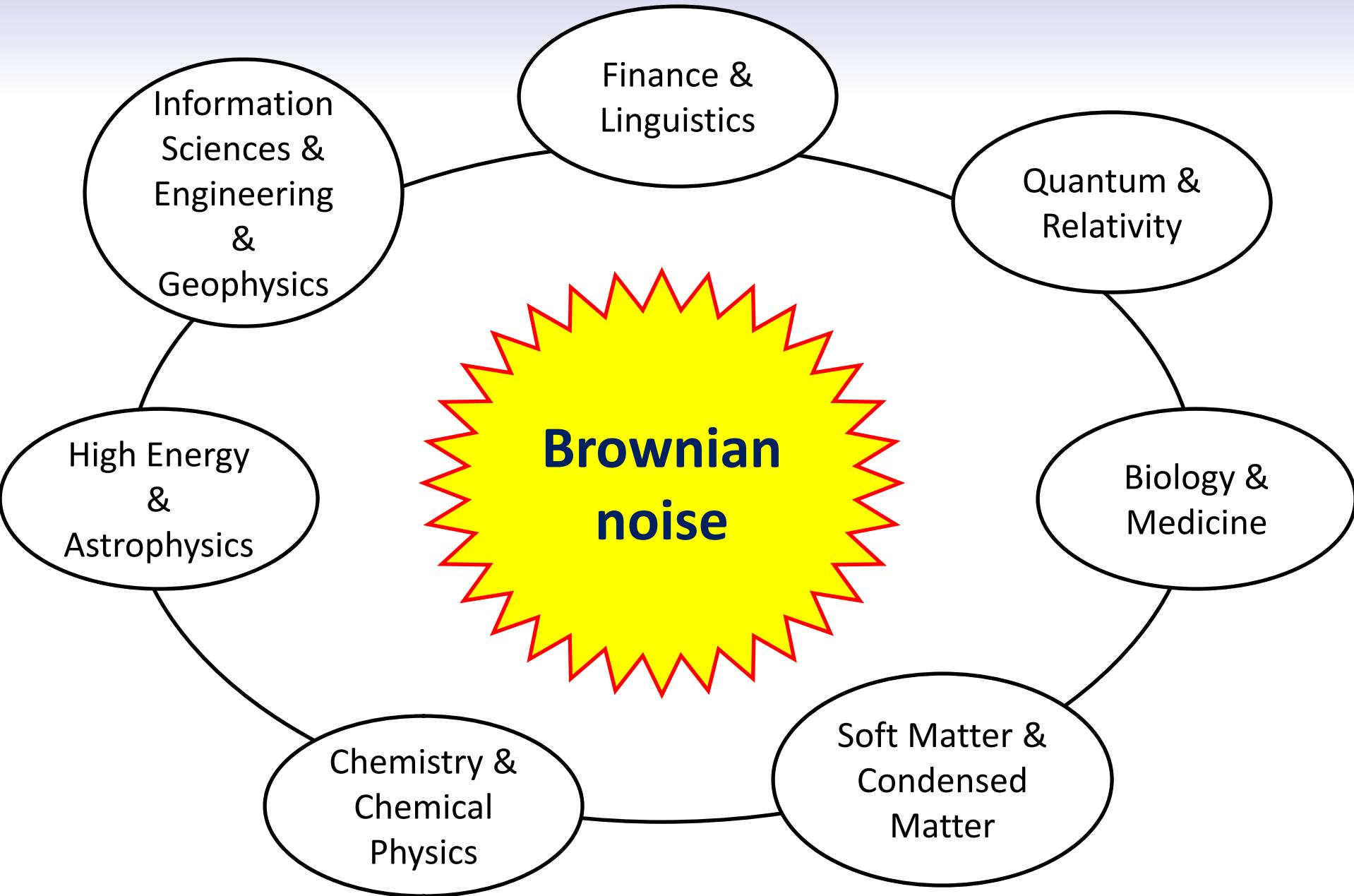
M. Smoluchowski (1872-1917)



Source: [wikipedia.org](https://en.wikipedia.org)

$$D = \frac{32}{243} \frac{mc^2}{\pi\mu R}$$

Ann. Phys. **21**, 756 (1906)



Quantum-Mechanics

= Brownian Motion ?

= Stochastic Mechanics ?

(E. Nelson; 1966, 1986)

$$p(x, t) = |\Psi(x, t)|^2$$



Schrödinger-
gleichung

$$\Psi(x, t) = |\Psi(x, t)| e^{iS(x, t)}$$

$$\dot{p}(x, t) = -\frac{\hbar}{m} \nabla [(\nabla \ln |\Psi(x, t)| + \nabla S(x, t)) p(x, t)] + \frac{\hbar}{2m} \nabla^2 p(x, t)$$

$$\mathbf{f}_1 \geq 0, \mathbf{f}_2 \geq 0 : \quad \frac{1}{2} \langle \mathbf{f}_1(t_1) \mathbf{f}_2(t_2) + \mathbf{f}_2(t_2) \mathbf{f}_1(t_1) \rangle \geq 0$$

QM: NO !

Diffusion: space-time only

classical diffusion (Markovian)

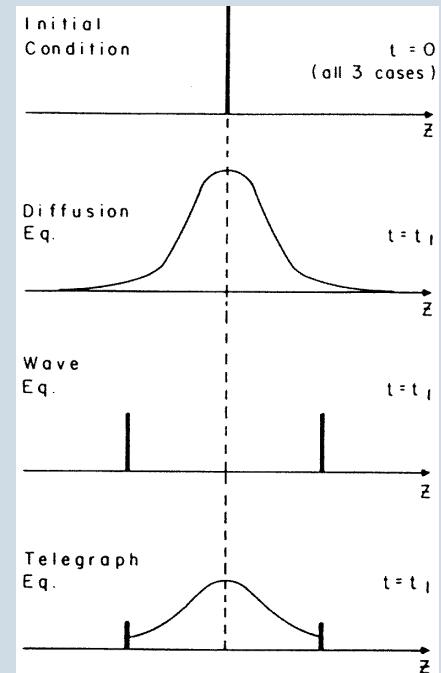
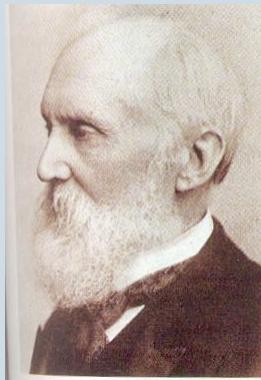
$$p(t, x|t_0, x_0) = \left[\frac{1}{4\pi \mathcal{D}(t - t_0)} \right]^{1/2} \exp \left[-\frac{(x - x_0)^2}{4\mathcal{D}(t - t_0)} \right].$$

telegraph equation (non-Markovian)

$$\tau_v \frac{\partial^2}{\partial t^2} \varrho + \frac{\partial}{\partial t} \varrho = \mathcal{D} \nabla^2 \varrho,$$

alternative approach

$$p(\bar{x}|\bar{x}_0) \propto \int_{a_-(\bar{x}|\bar{x}_0)}^{a_+(\bar{x}|\bar{x}_0)} da \exp\left(-\frac{a}{2\mathcal{D}}\right)$$



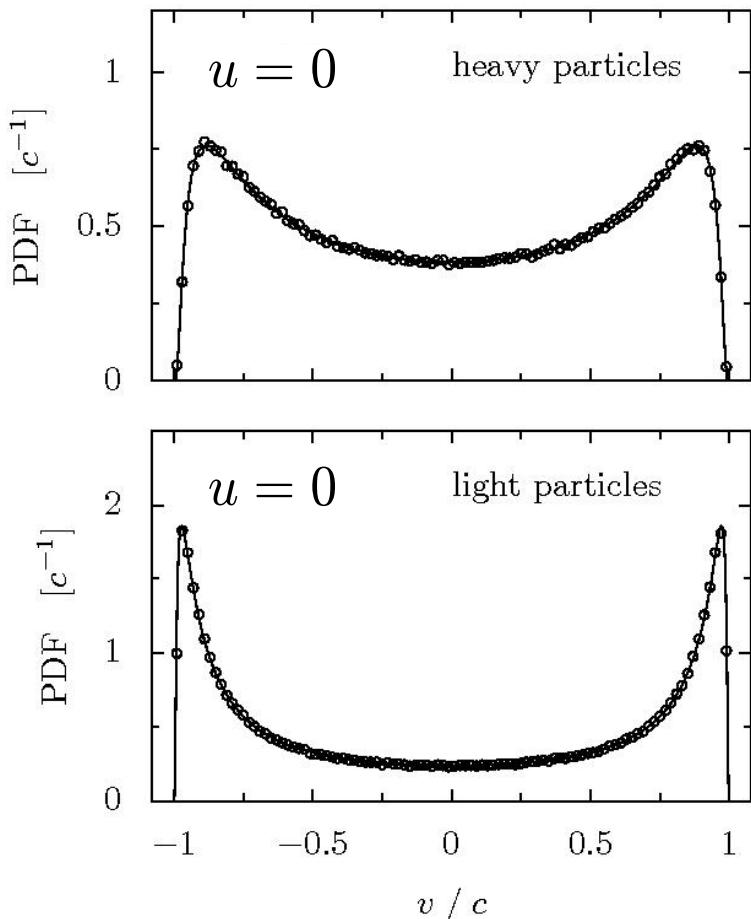
J Masoliver & G H Weiss
Eur J Phys 17:190 (1996)

PRD 75:043001 (2007)

Jüttner Gas

$$f_{\text{Maxwell}}(\vec{p}) = [\beta/(2\pi m)]^{d/2} \exp(-\beta p^2/2m)$$

$$f_{\text{Jüttner}}(\vec{p}) = Z_d^{-1} \exp\left[-\beta_J(m^2 c^4 + p^2 c^2)^{1/2}\right]$$



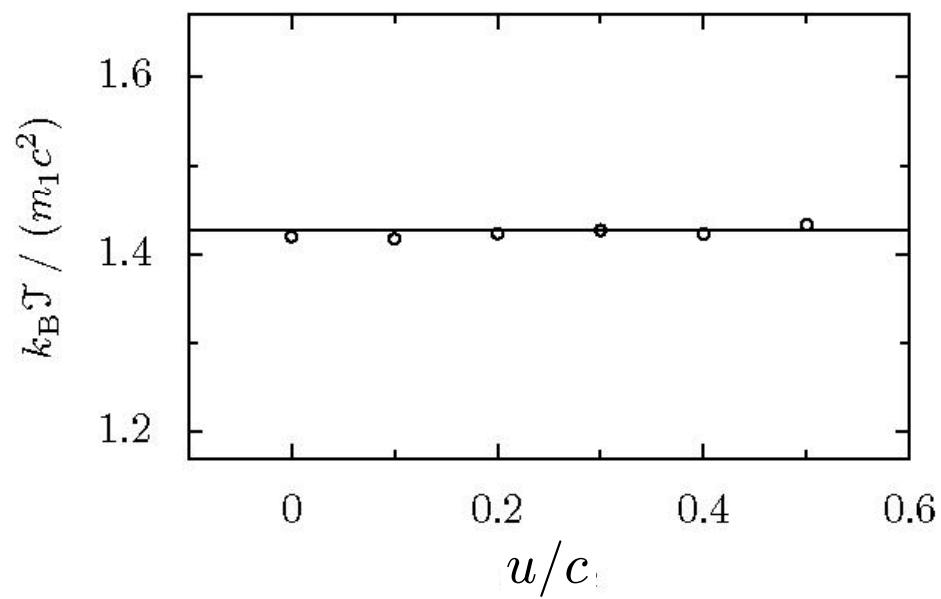
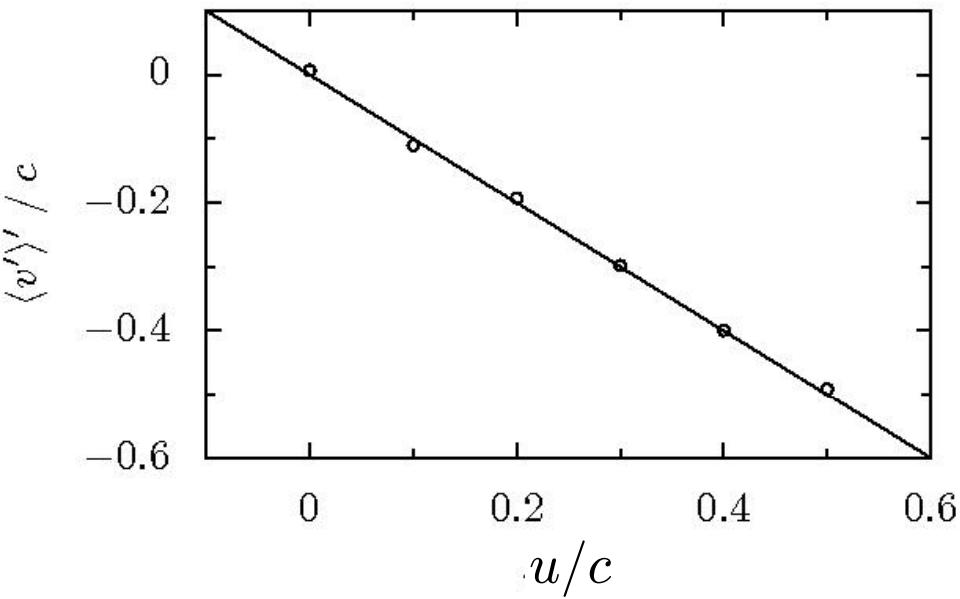
$\langle \vec{p} \cdot \vec{v} \rangle = dk_B \mathcal{T} = d/\beta_J$
statistical relativistic temperature
 $T = \mathcal{T} = (k_B \beta_J)^{-1}$

Measuring temperature in Lorentz invariant way

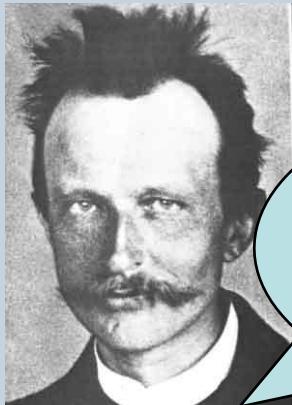
Lorentz invariant equipartition theorem

$$k_B T = m \gamma(u)^3 \langle \gamma(v') (v' + u)^2 \rangle_{t'}$$

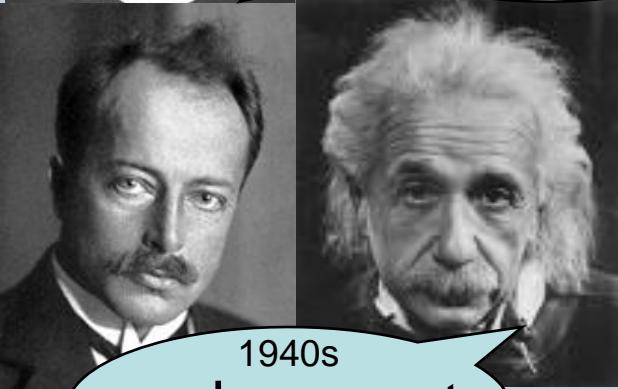
$$\text{with } u = -\langle v' \rangle_{t'}$$



“Temperature” problem in RTD?



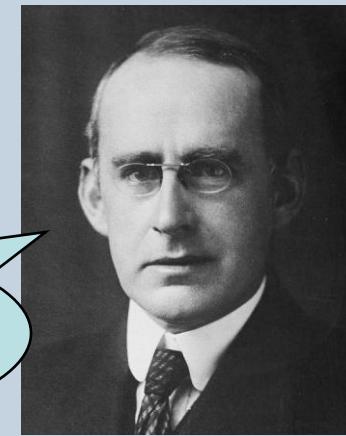
1907/08
moving bodies
appear cooler



1940s
maybe ... not

1923/1963

.. hotter!



$$T'(w) = T (1 - w^2)^{\alpha/2}$$

$$\alpha = \begin{cases} +1 & \text{Planck, Einstein} \\ 0 & \text{Landsberg, van Kampen} \\ -1 & \text{Ott} \end{cases}$$

$T=T'$ 1966-69

CK Yuen, Amer. J. Phys. 38:246 (1970)



Two prominent examples

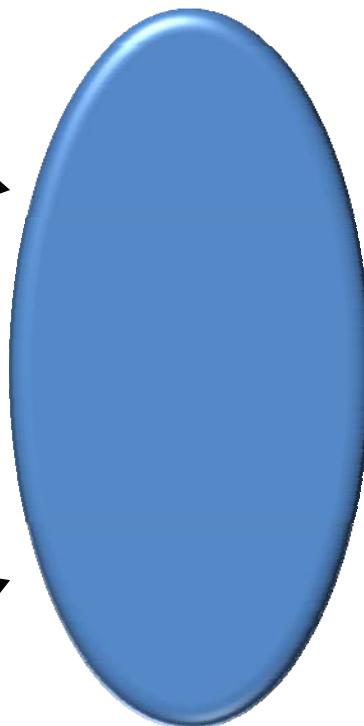
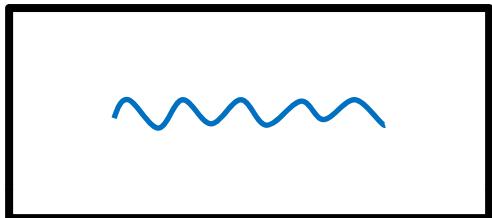
Stochastic
Resonance

Brownian
Motors

Stochastic Resonance

(in a nutshell)

Weak signal

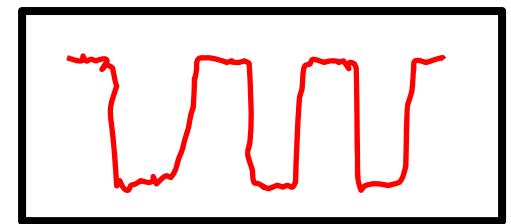


Noise source



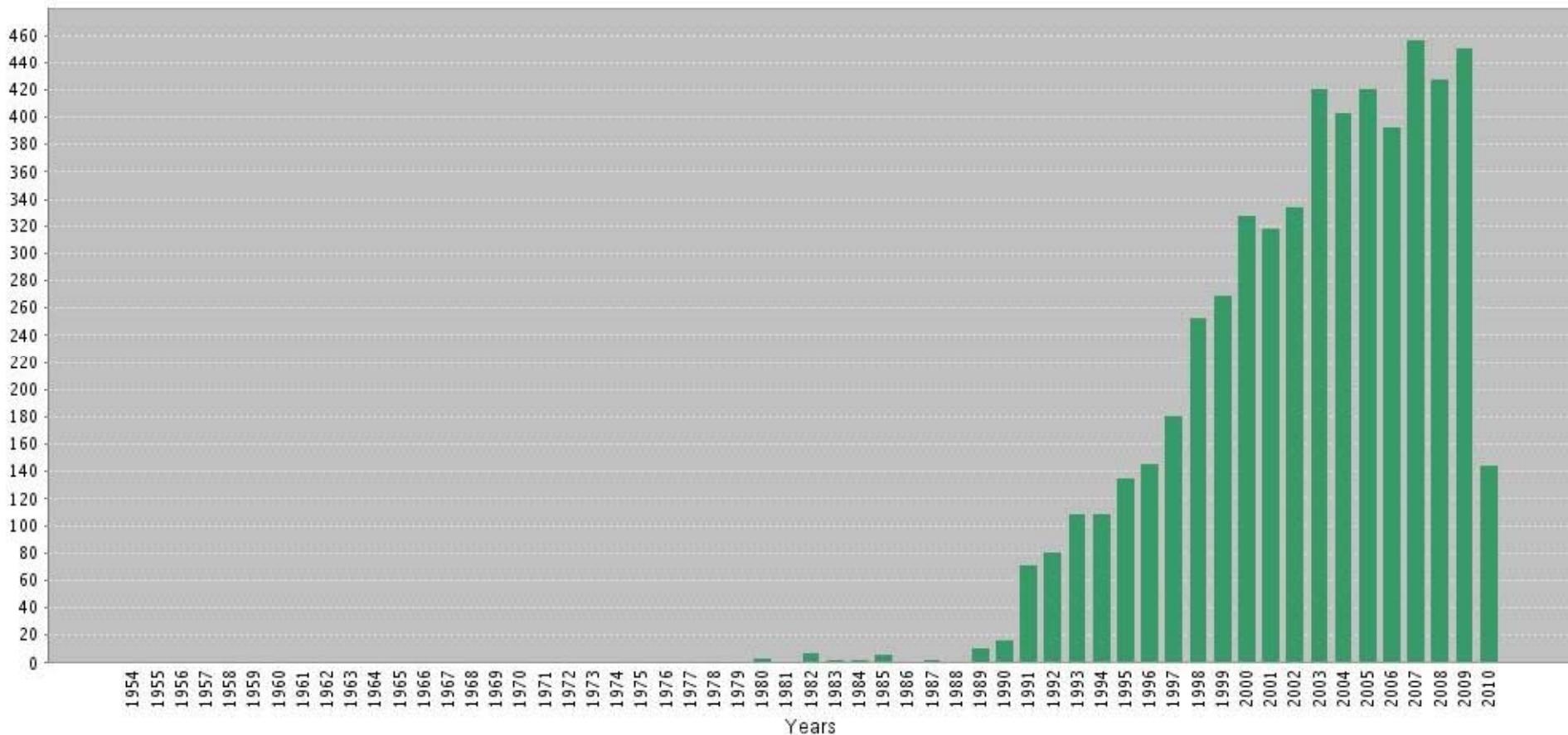
System

Output signal



SR - Citations

Published Items in Each Year

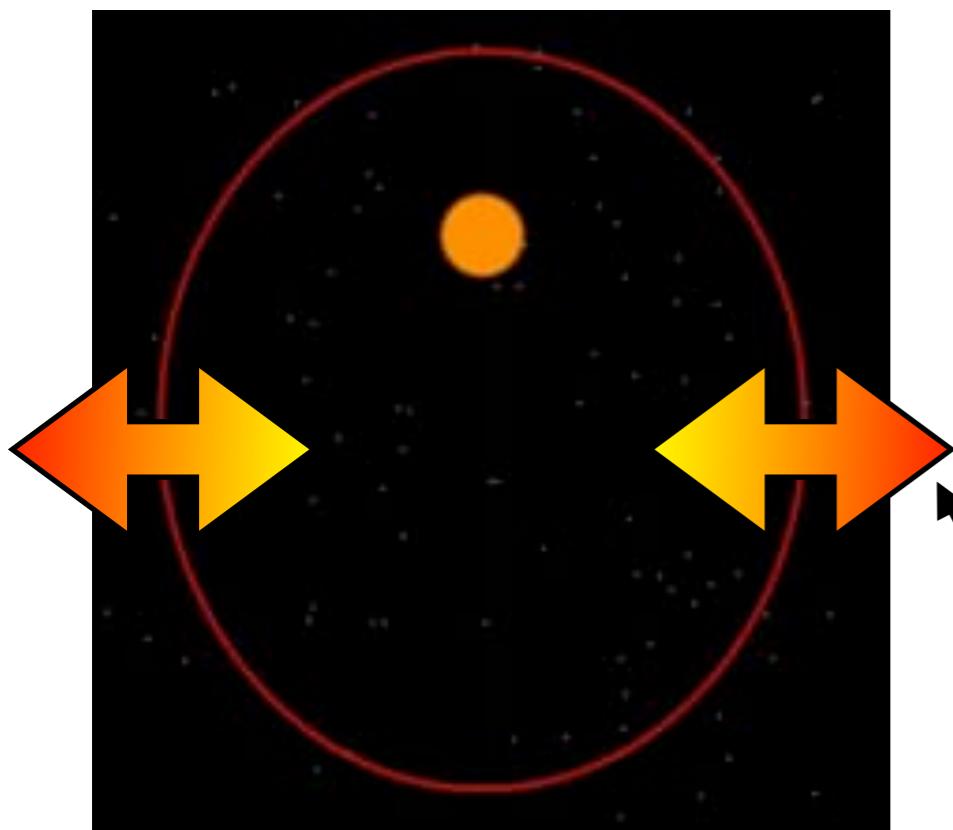


papers in 2009: ≈ 460
> 85000 cites in total

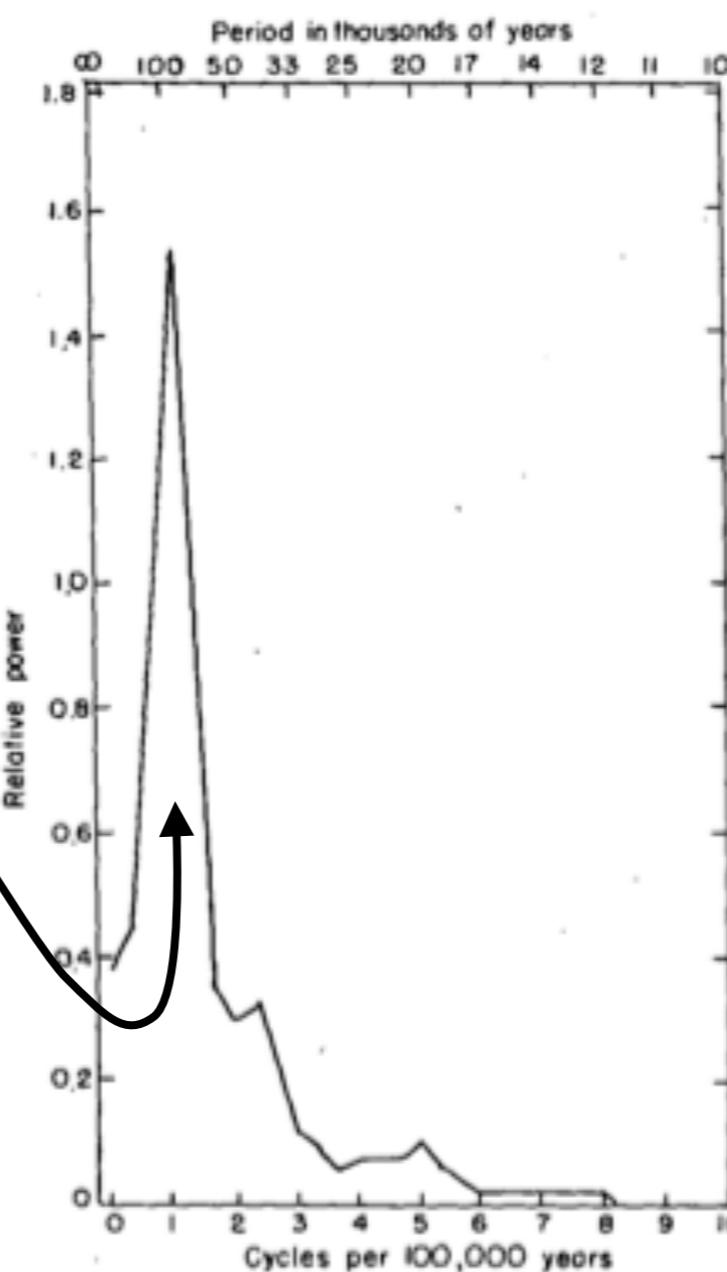
Why are the ice-ages so periodic ?

Milankowitch cycles:

Small changes in earth orbit eccentricity with 100k year periodicity



M. Milankowitch, Handbuch der Klimatologie I
(1930)

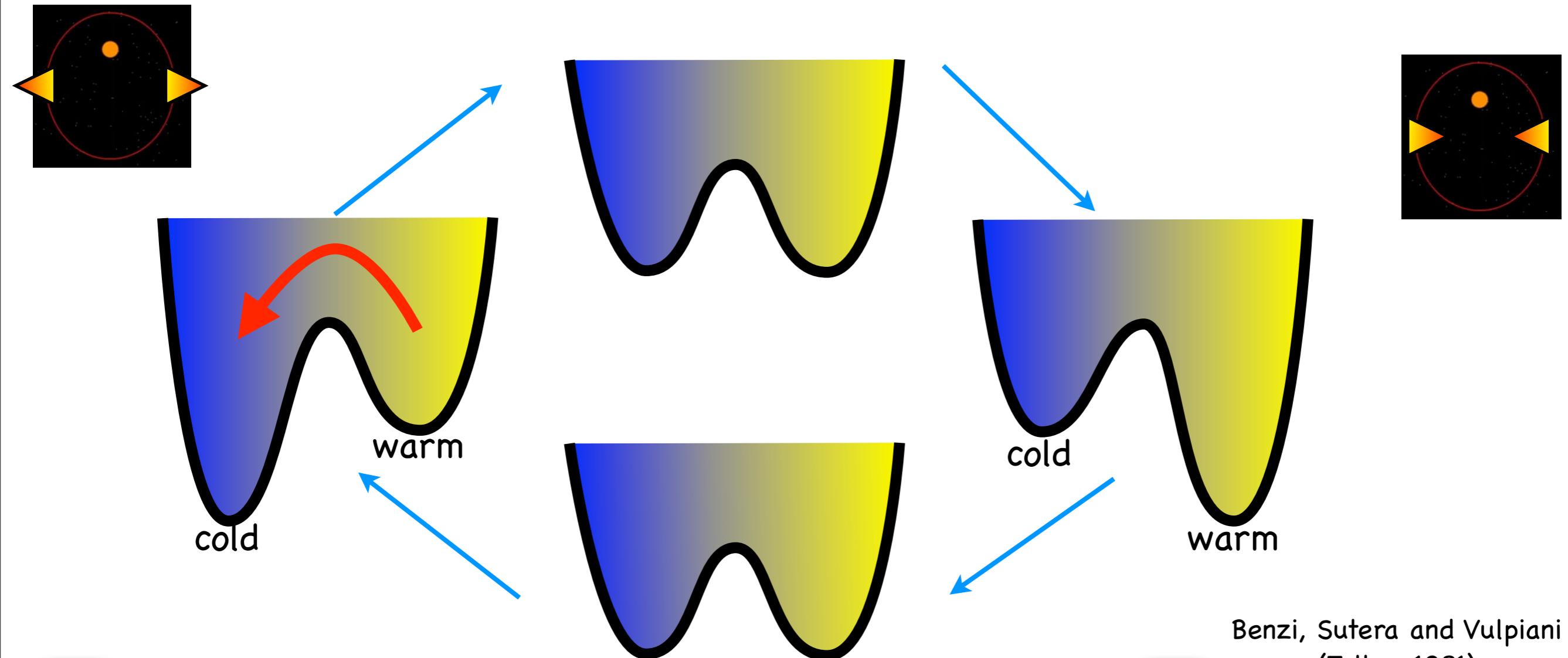


Changes are small!
(<0.1% of solar
constant)

What can amplify
those small changes ?

Milankowitch Cycles and Bistability

Climate "landscape"

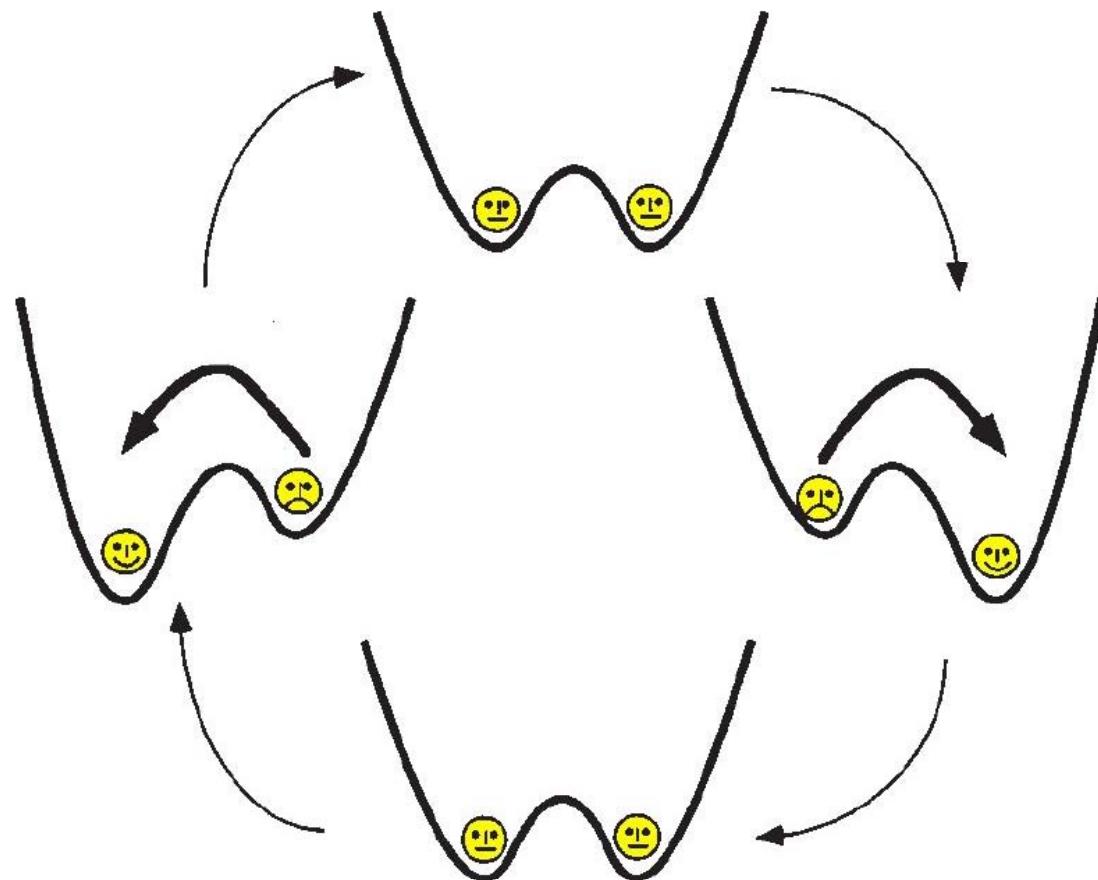


Benzi, Sutera and Vulpiani
(Tellus, 1981)

C. Nicolis and G. Nicoli
(Tellus, 1981)

- The 100ky cycles only bias the climate
- Fluctuations make climate switch
- small changes of conditions can have huge impact

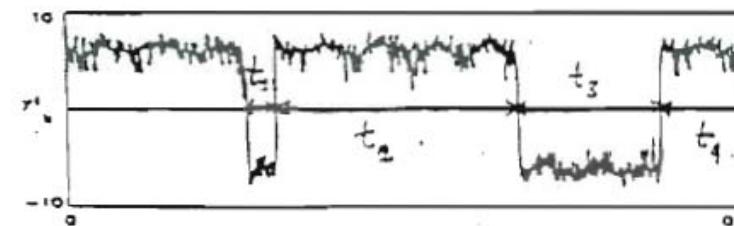
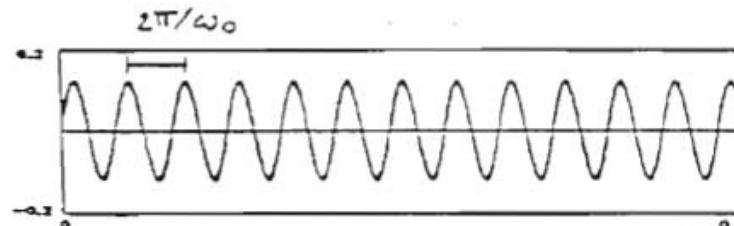
Noise-assisted synchronized hopping



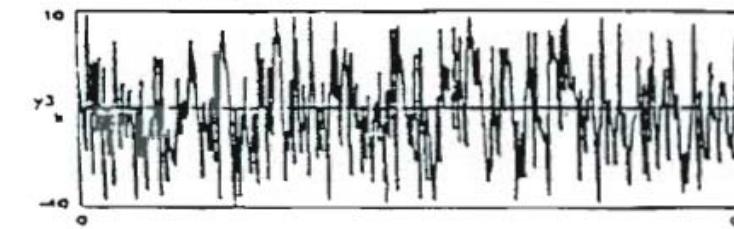
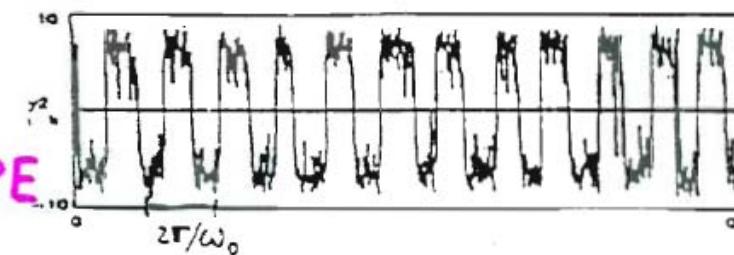
$$T_{\text{period}} \simeq 2T_{\text{escape}}$$

Synchronization

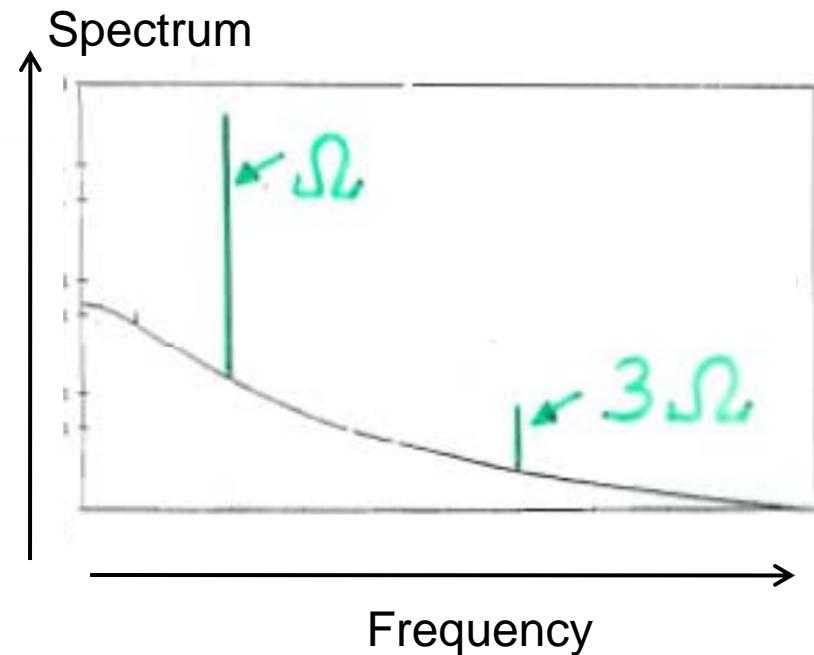
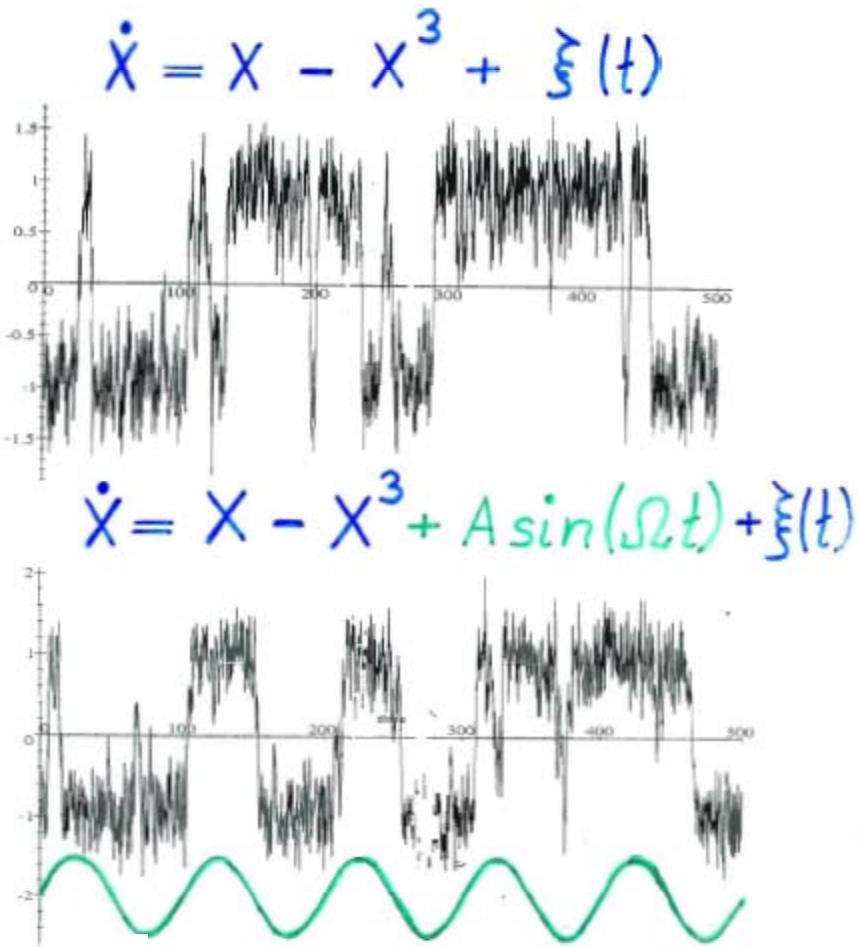
SIGNAL



$T_e \sim 2\Gamma^{-1}$
ESCAPE



Power spectral density



Measuring SR

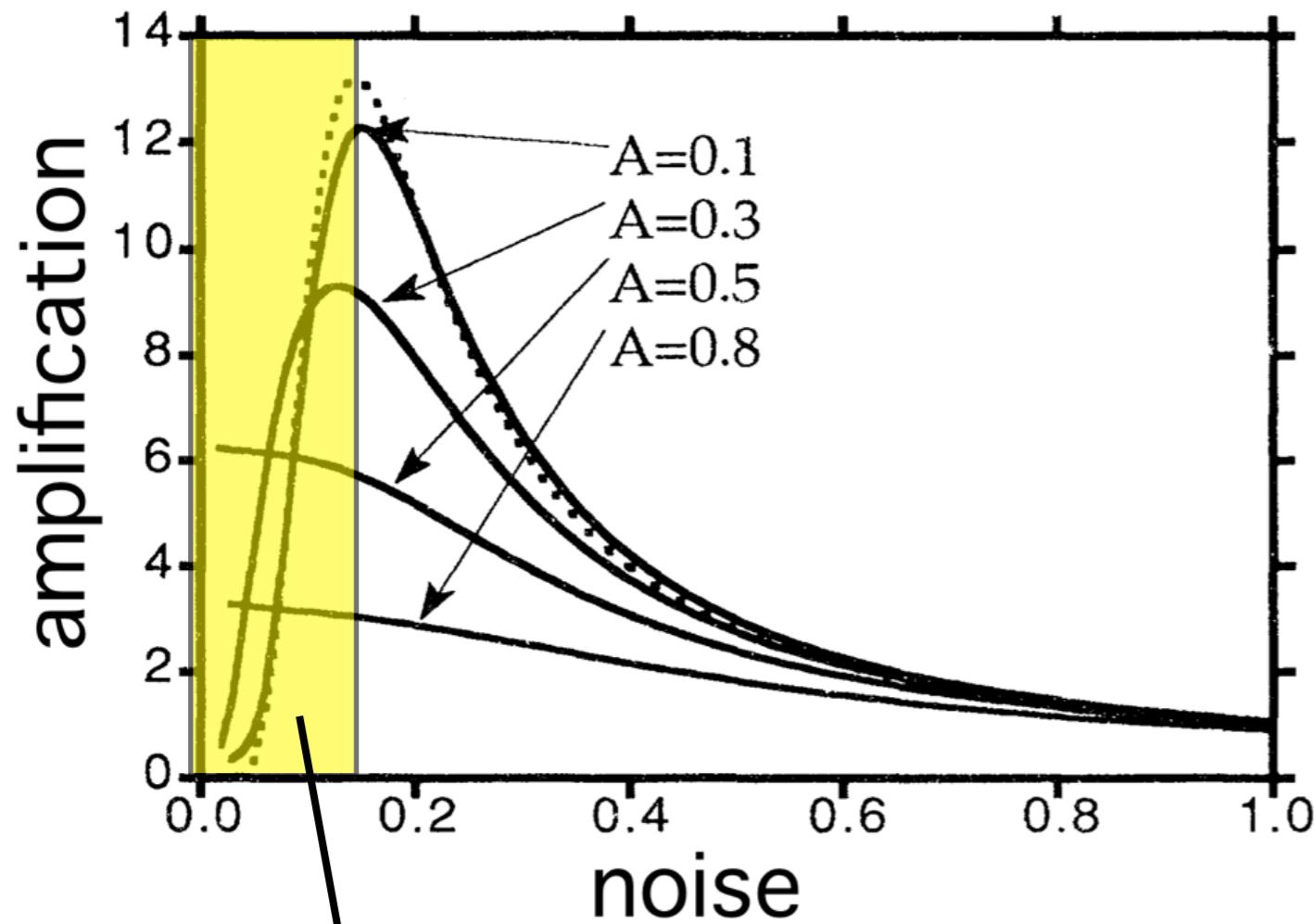
- Signal to noise ratio
- Spectral amplification
- mutual information
- cross-correlation: input \leftrightarrow output
- peak area, (phase-) synchronization, ...

SR-reviews:

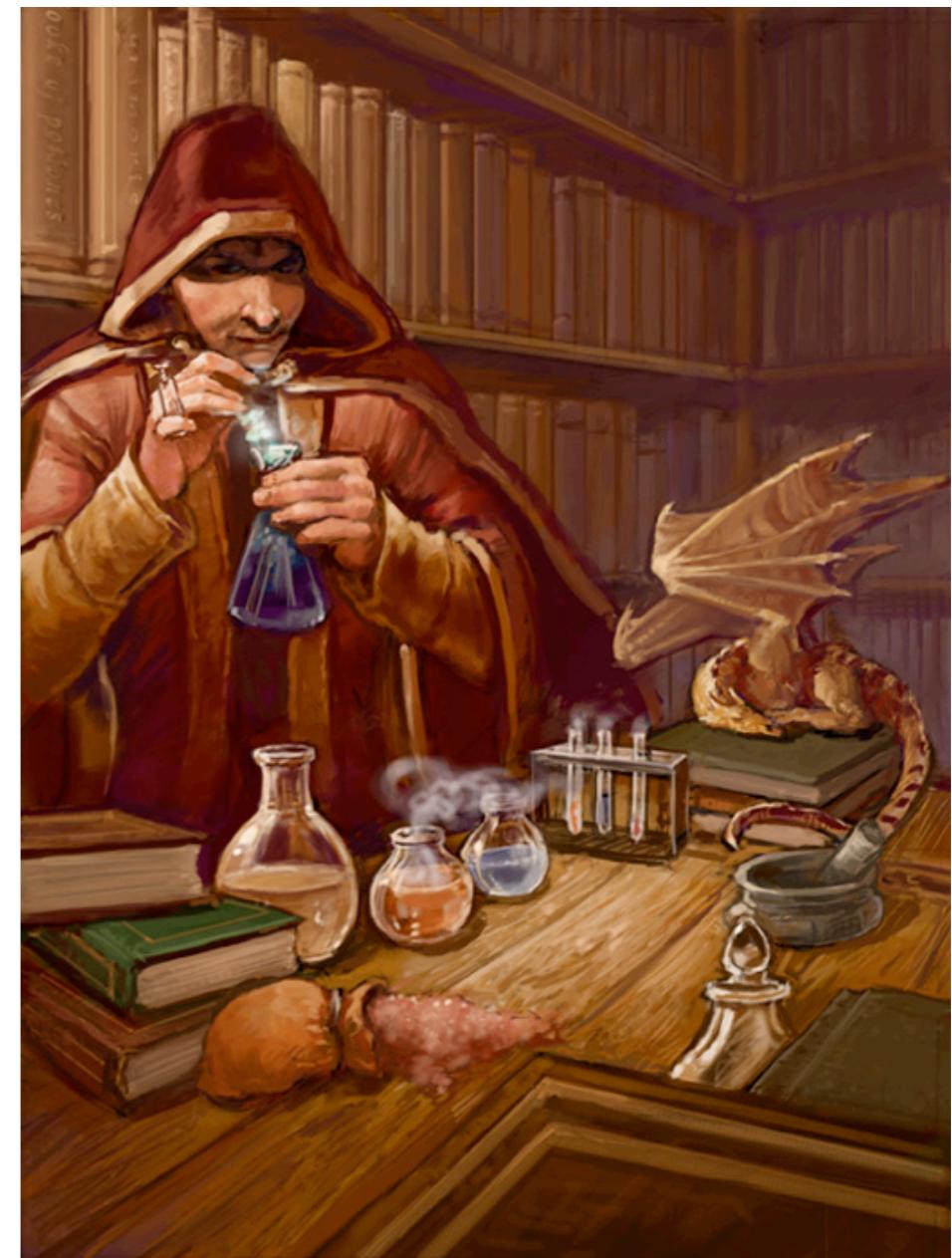
L. Gamaitoni, P. Hänggi, P. Jung, F. Marchesoni, Rev. Mod. Phys. **70**, 223 (1998)
P. Hänggi, ChemPhysChem **3**, 285 (2002)

Amplification of small signals by noise

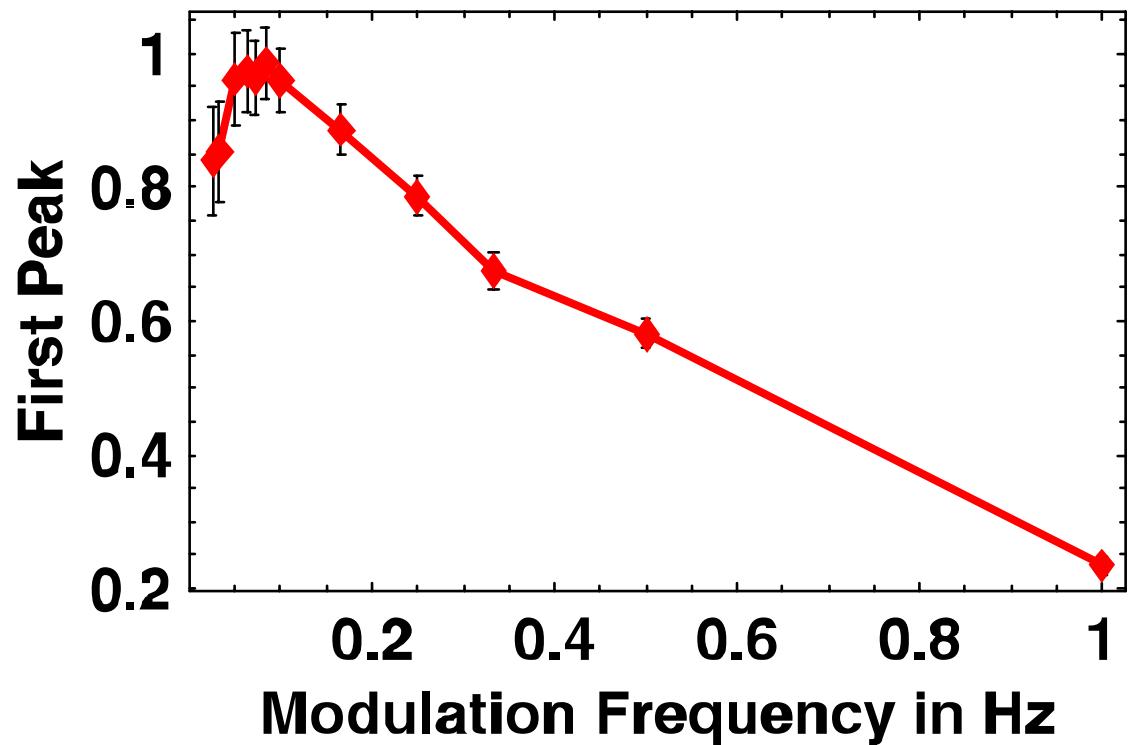
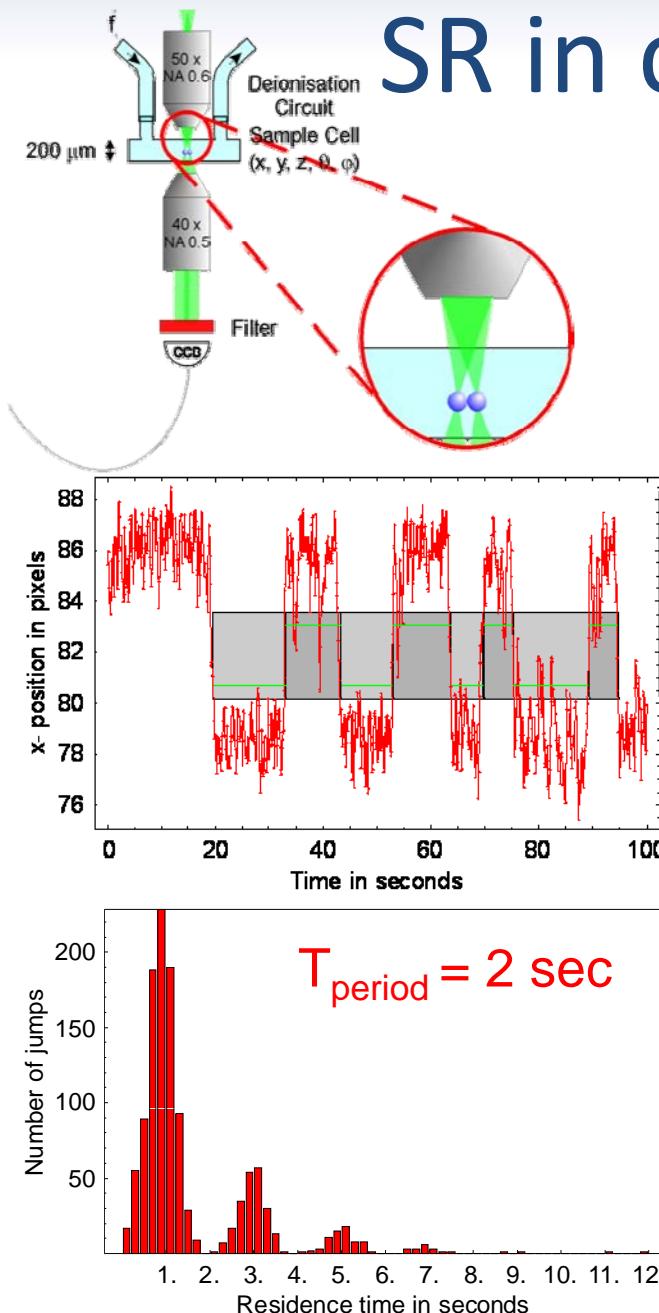
(P. Jung, P. Hänggi, Phys. Rev. A **44**, 8032 (1991))



More noise , more signal !!



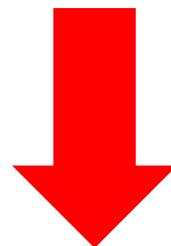
SR in colloidal systems



D. Babic, C. Schmitt, I. Poberaj, C. Bechinger,
Europhys. Lett. **67**, 158 (2004)

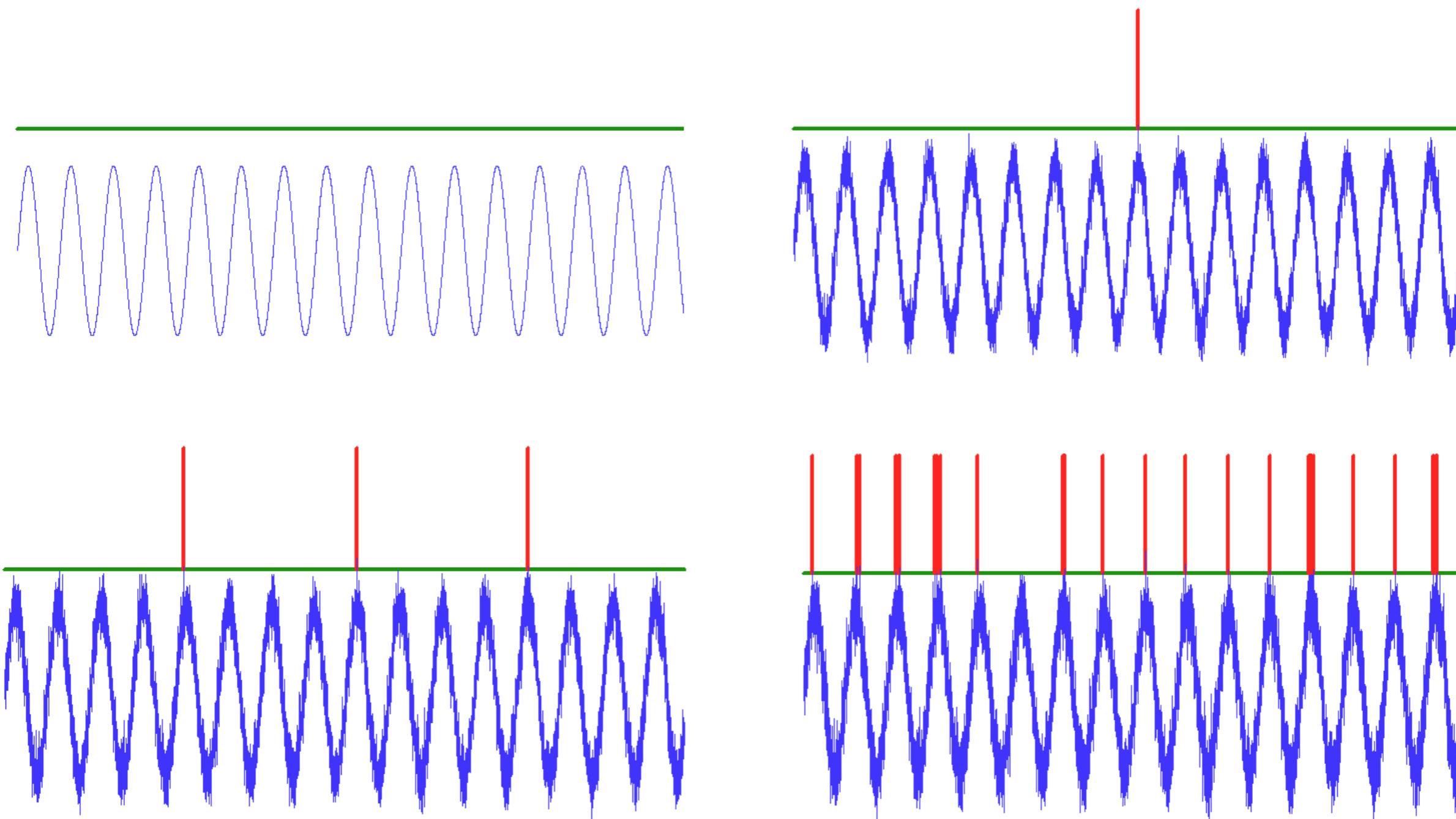
SR - Ingredients

- ✓ Threshold system
- ✓ Weak (subthreshold) signal
- ✓ Noise



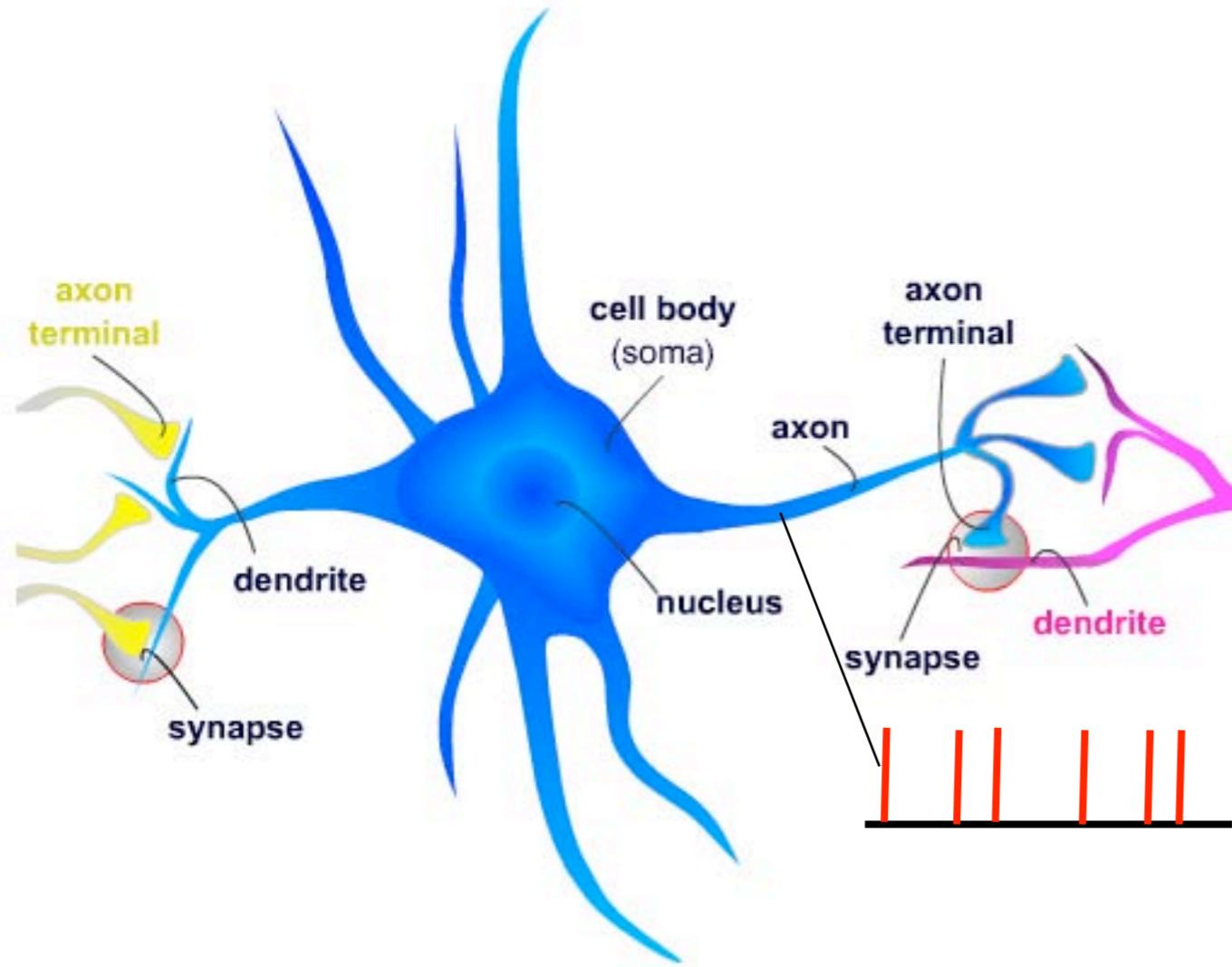
Anomalous amplification properties

Thresholds and Stochastic Resonance



P. Jung, Phys. Rev. E50, 2513 (1994), F. Moss and L. Kiss, EPL, 29 (1995)

Stochastic Resonance in Neurobiology



Input: currents at synapses

Processing: action potential if the sum of currents exceeds threshold

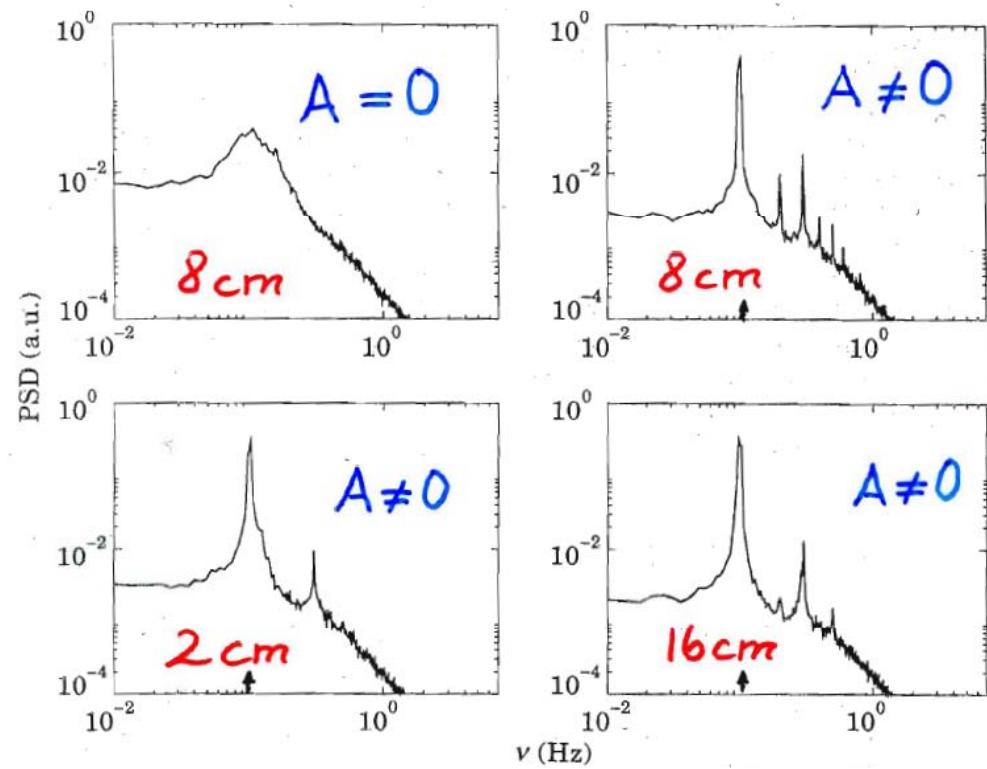
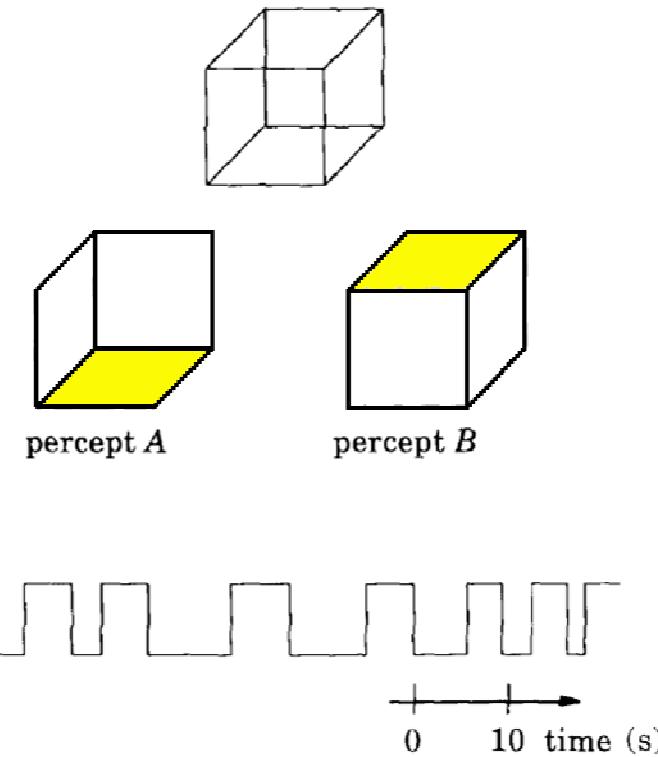
Output: electric pulses traveling down the axon

source: Consortium on Cognitive Science Instruction (CCSI)

Basic idea: Signals below threshold can be detected in the presence of additional noise

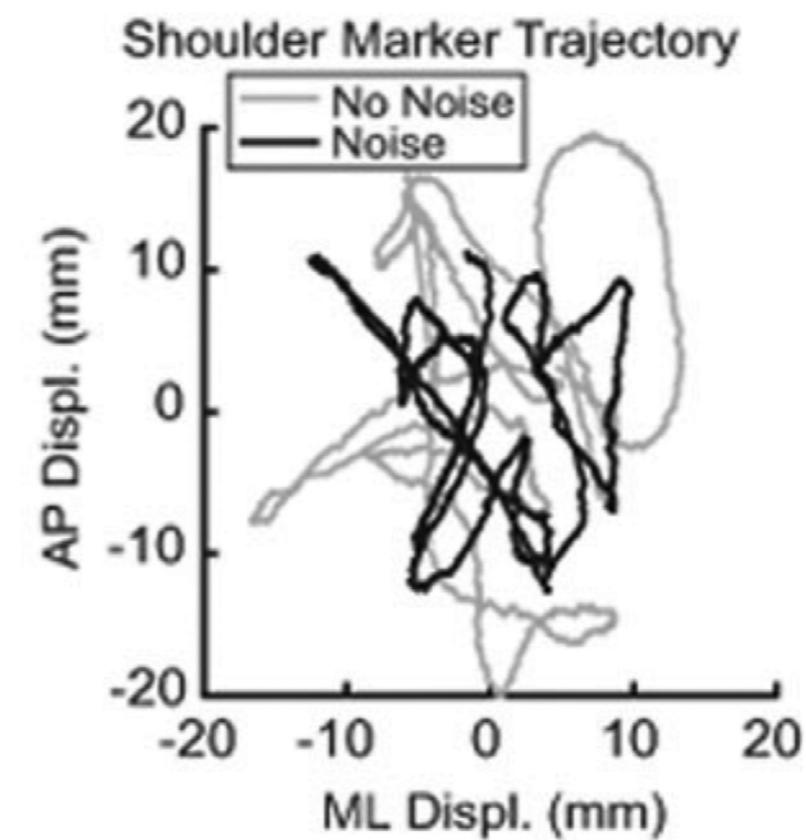
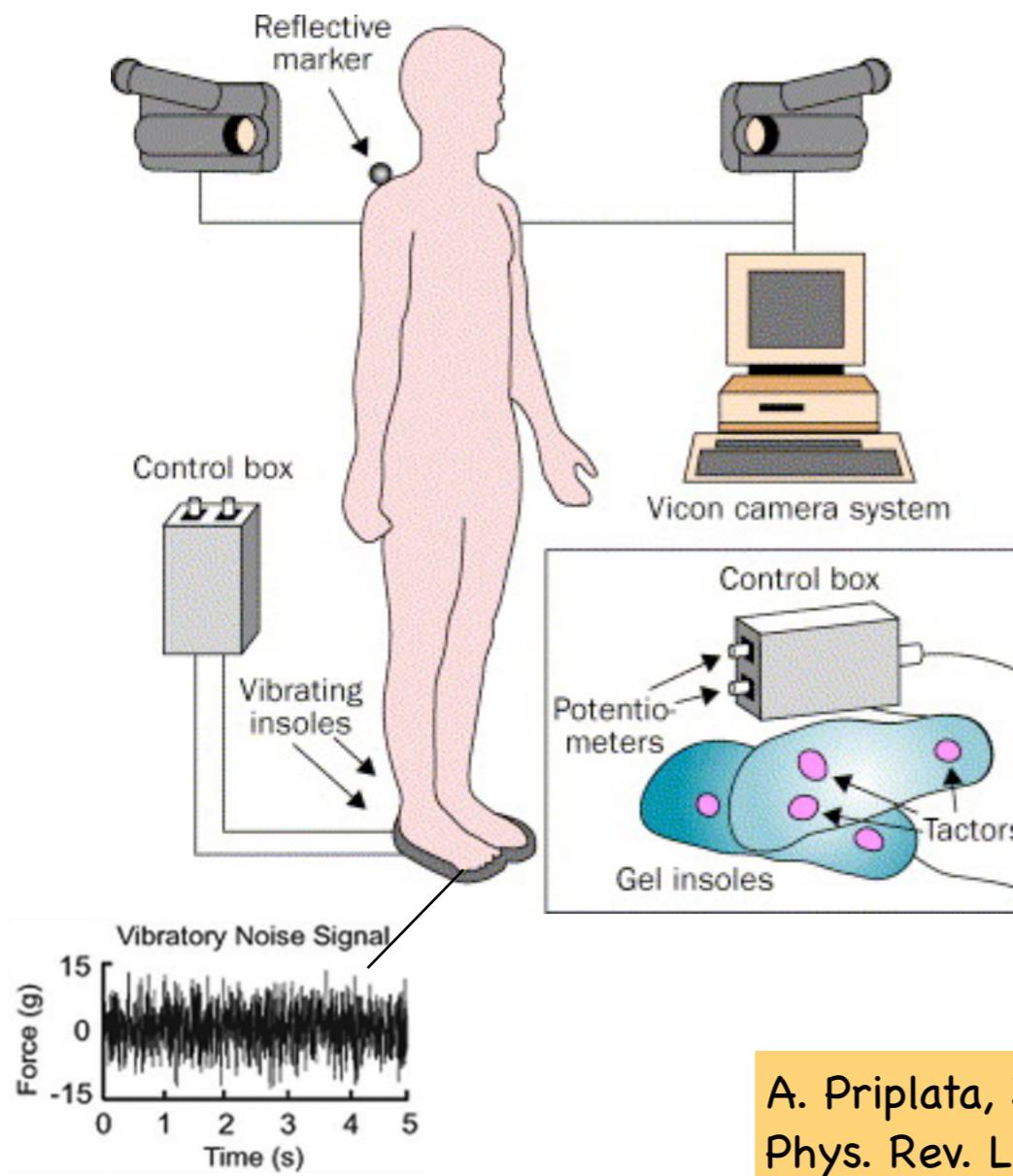
SR in Visual Perception

Experiment:



SR and human posture control

Somatosensory function declines with age and in diabetic patients. Can additional noise help restore function?



Reduction in sway of person

A. Priplata, J. Niemi, M. Salen, J. Harry, L.A. Lipsitz and J.J. Collins
Phys. Rev. Lett. 89 (2002)



FREI
AKTIVE REHA-SYSTEME

THERAPIE.
PRÄVENTION.
TRAINING.

srt-zeptoring®
und
srt zeptor®

- Parkinson
- Multiple sclerosis (MS)
- Stroke / skull-brain-trauma
- Cross-section paralysis
- Depression
- Pain
- ...

[SRT Zeptor Training - Powerslide Team](#)



FREI
AKTIVE REHA-SYSTEME

THERAPIE.
PRÄVENTION.
TRAINING.

srt-zeptoring®
und
srt zeptor®

- Parkinson
- Multiple sclerosis (MS)
- Stroke / skull-brain-trauma
- Cross-section paralysis
- Depression
- Pain
- ...

[SRT Zeptor Training - Powerslide Team](#)

SR trends

- Spatio – temporal SR
- Aperiodic SR
- Quantum SR

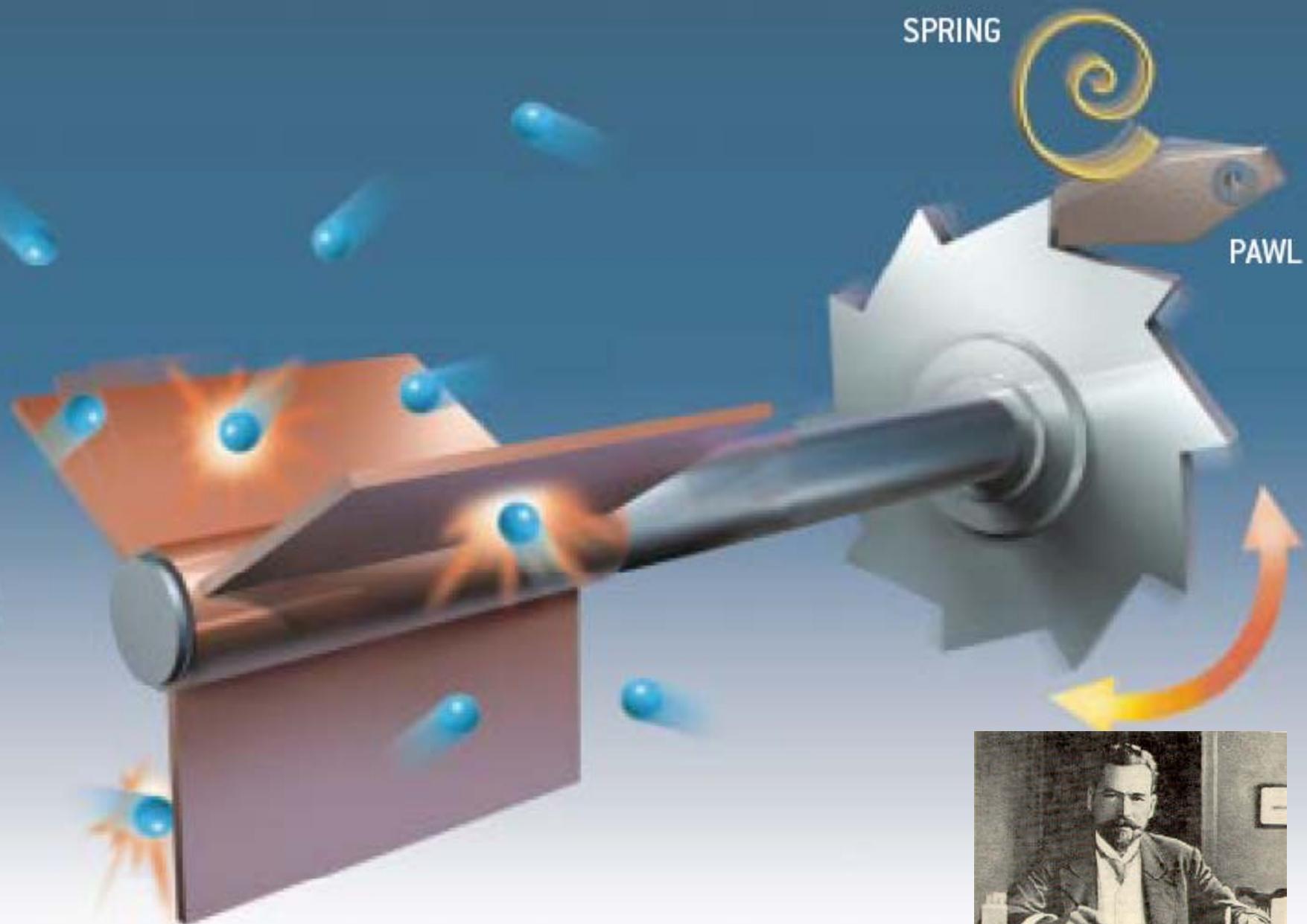
Motors \implies Brownian motors

Two heat reservoirs

One heat reservoir

Perpetuum mobile of the second kind?

NO !



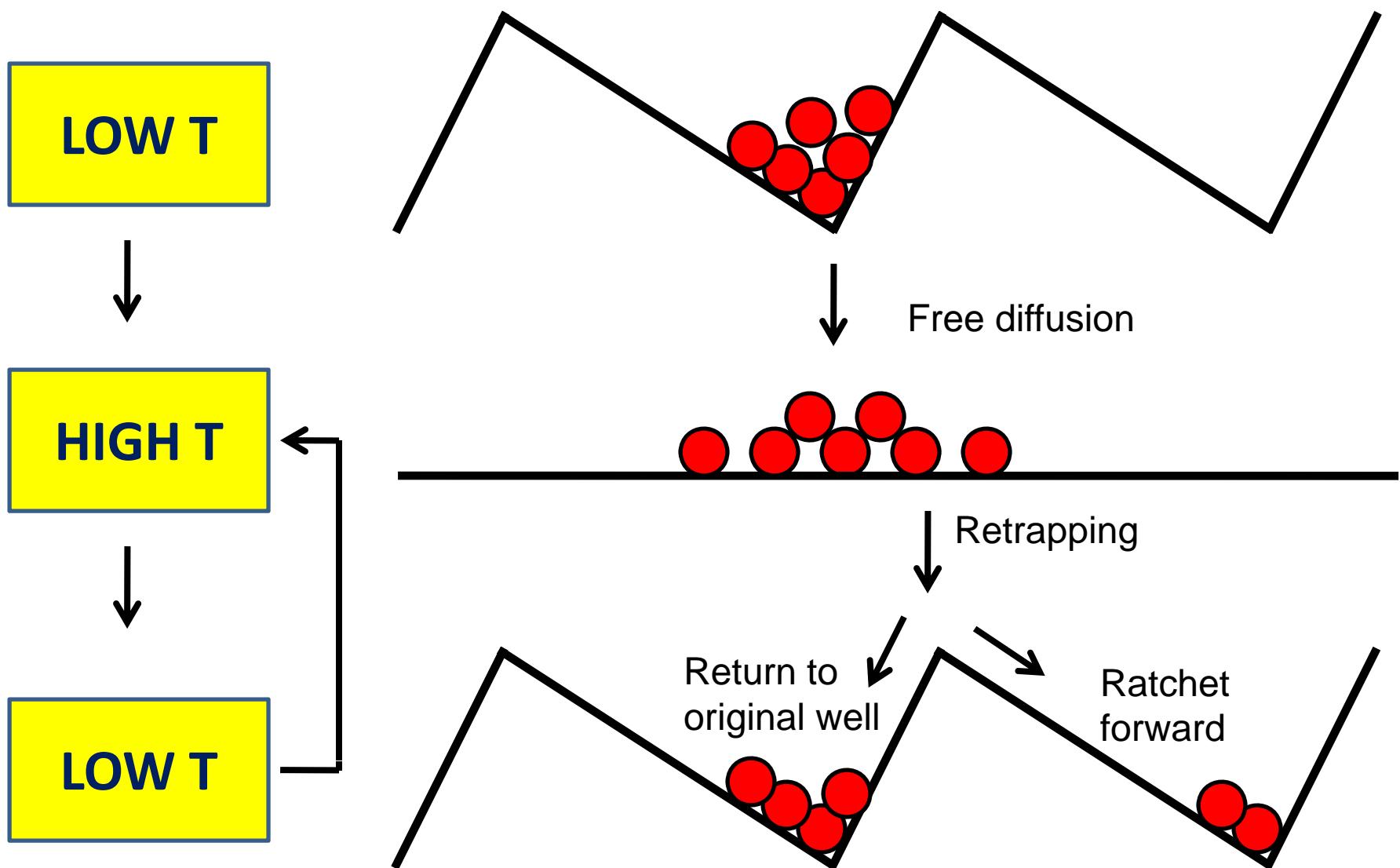
Source: Scientific American (2001)



Brownian motor

Movie

Temperature / Flashing Ratchet



Brownian motors - Characteristics

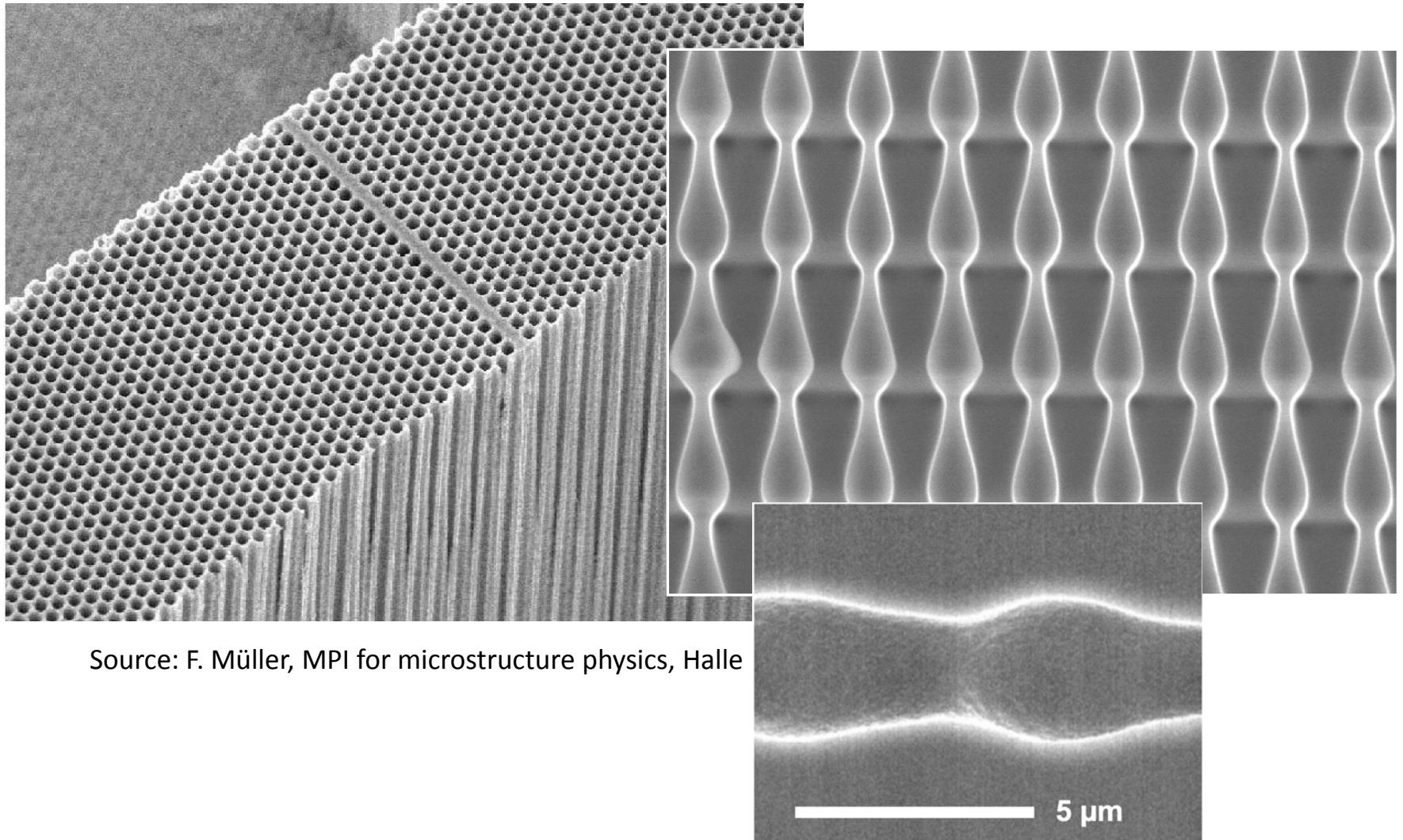
- Noise & AC-Input → **DC-Ouput**
- Non-equilibrium Noise → **Directed Transport**
- **Current reversals**
- **Applications:**
 - Novel pumps and traps for charged or neutral particles
 - Brownian diodes & transistors

Ask not what physics can do
for biology, ask what biology
can do for physics

REVIEWS OF MODERN PHYSICS, VOLUME 81, JANUARY–MARCH 2009

Artificial Brownian motors: Controlling transport on the nanoscale
P.H. and F. Marchesoni

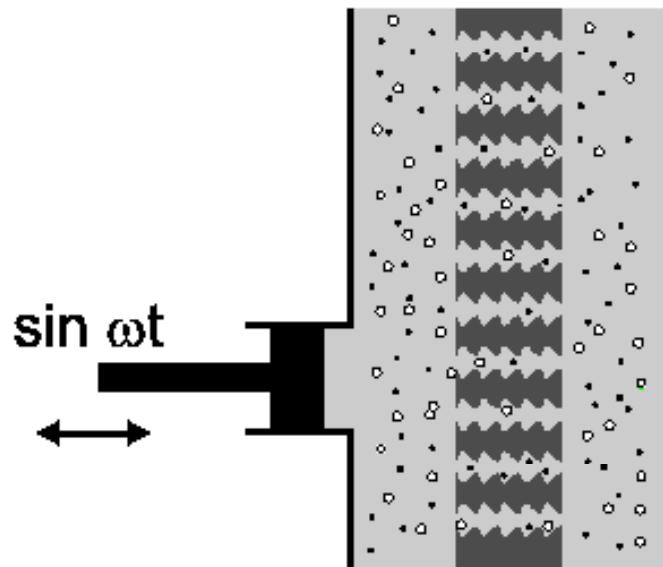
Drift Ratchet - Device



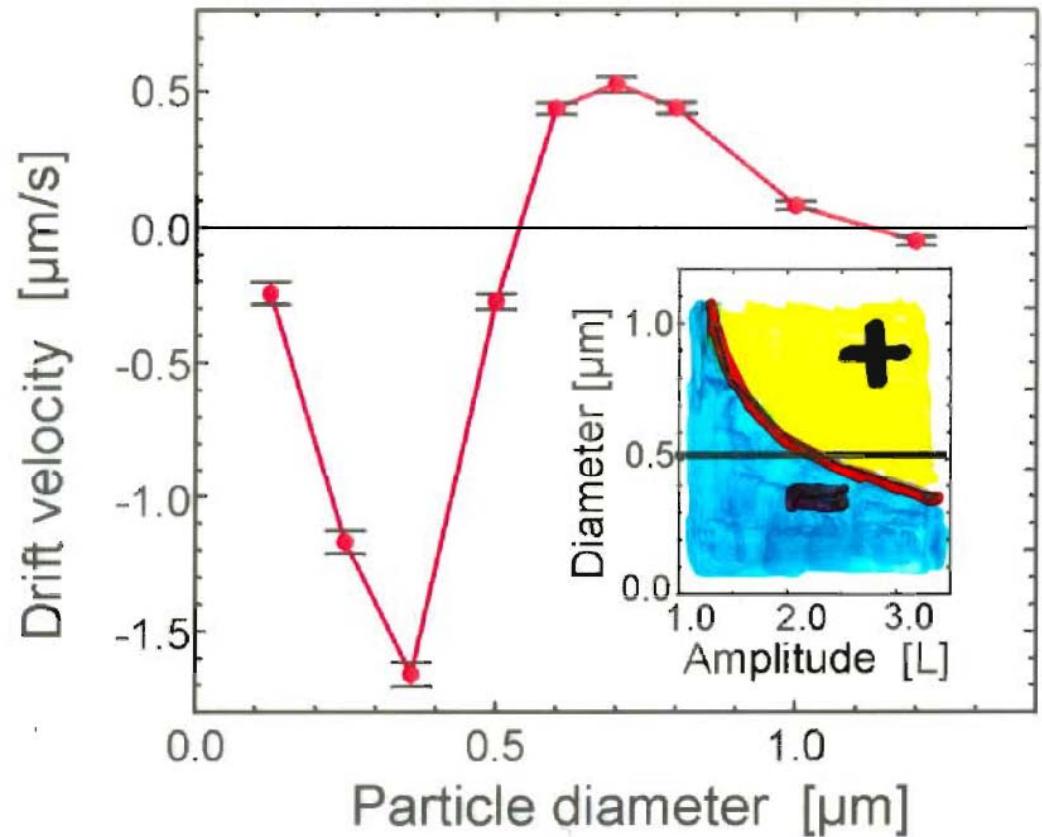
Drift Ratchet - Theory

C. Kettner, P. Reimann, P. H., F. Müller, Phys. Rev. E **61**, 312 (2000)

Setup

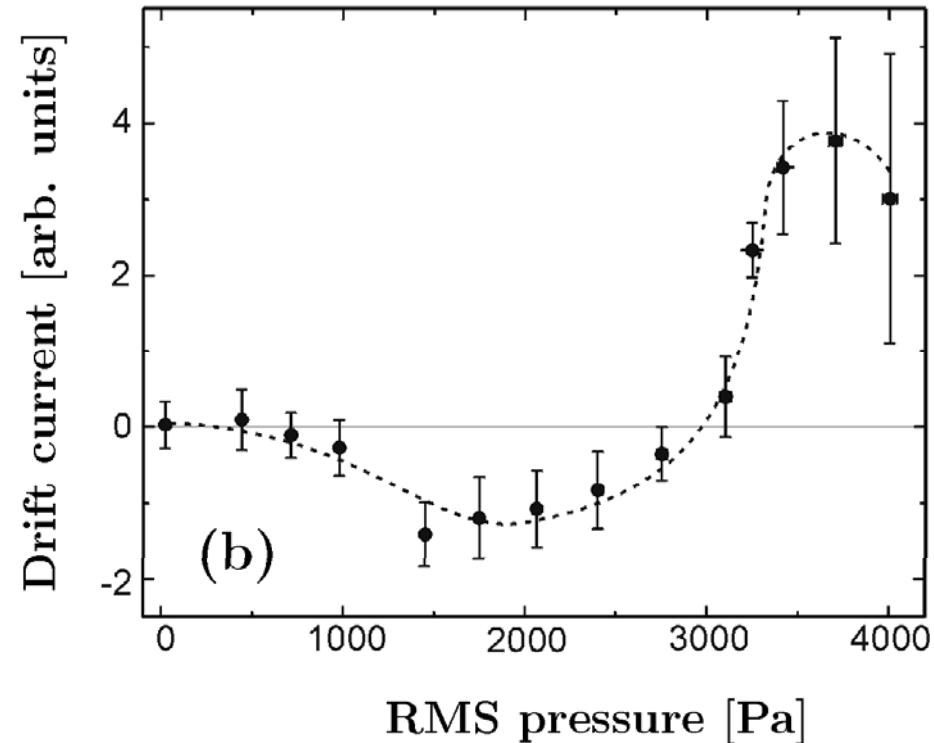
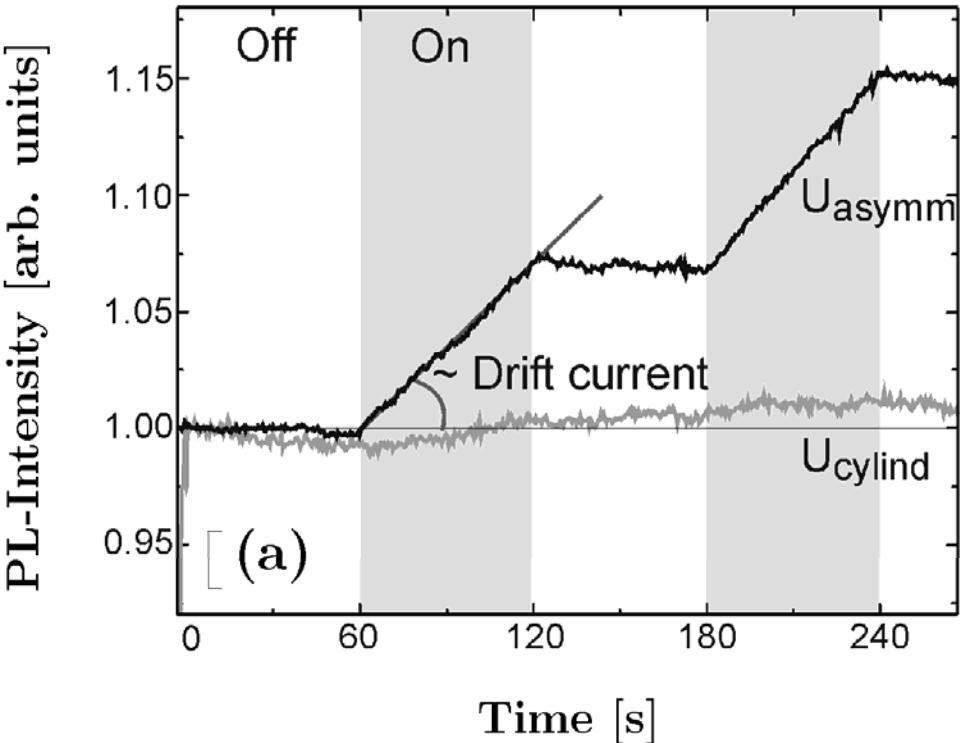


Particle Separation



Drift Ratchet – Experiment

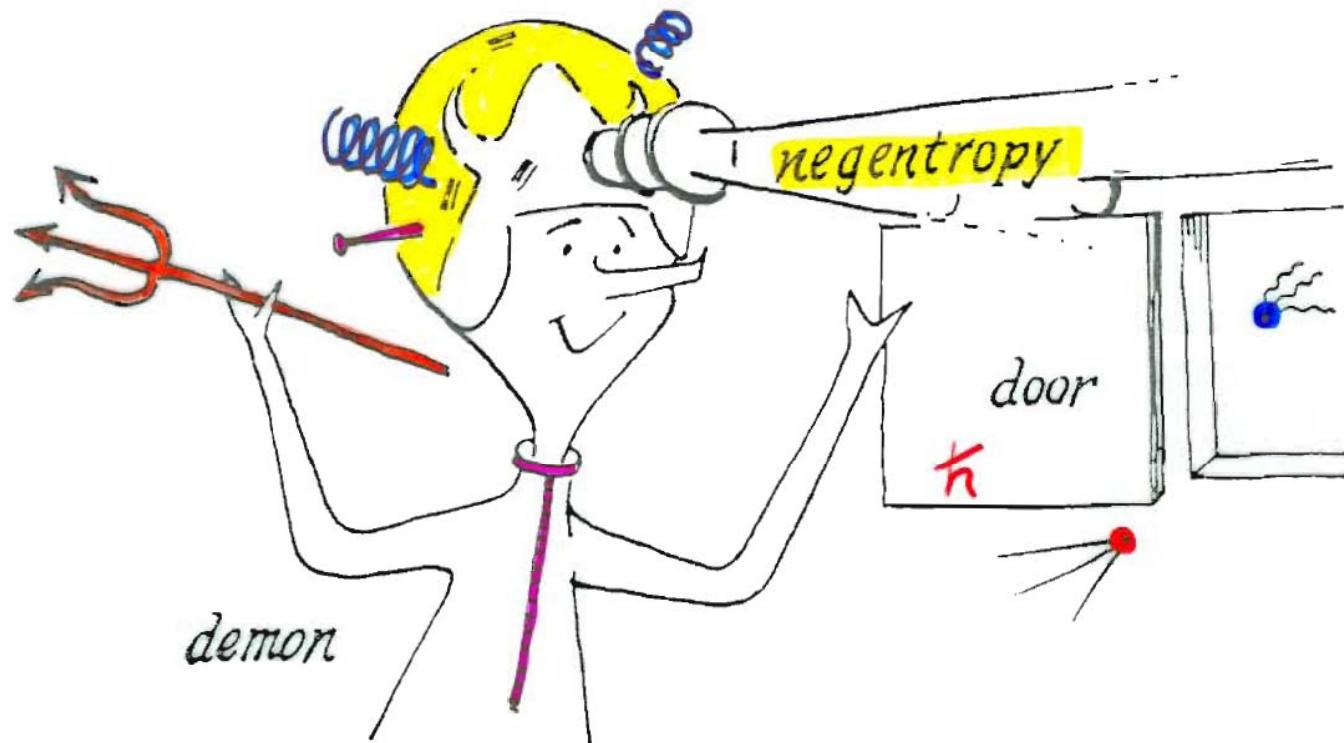
S. Matthias, F. Müller, Nature 424, 53 (2003)





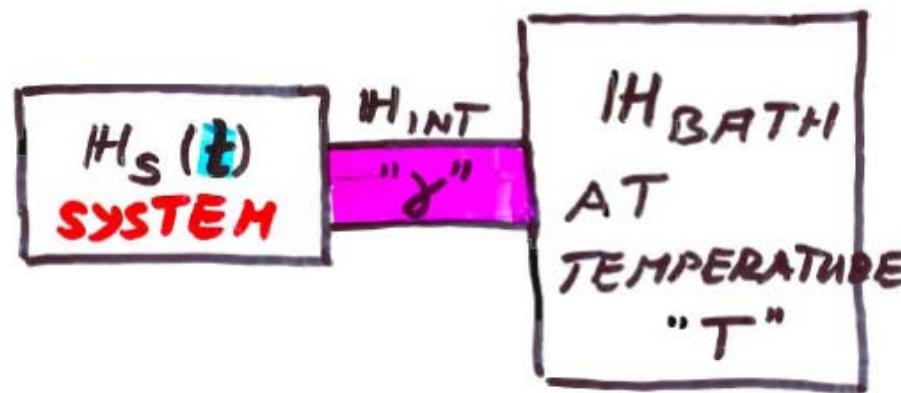
Quantum Demon ?

A measurement → Increase information → Reduction of entropy



Source: H.S. Leff, *Maxwell's Demon* (Adam Hilger, Bristol, 1990)

Quantum Brownian Motors



$$i\hbar \dot{\varrho} = [H_S(t) + H_{INT} + H_{BATH}, \varrho]$$

Hilbert space: $SYSTEM \otimes BATH$

SUPER-
BATH

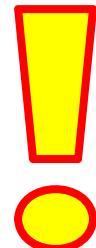
Quantum-Langevin-equation

$$m\ddot{\mathbf{x}}(t) + \int_{-\infty}^t \gamma(t-t') \dot{\mathbf{x}}(t') dt' + V'(\mathbf{x}; t) = \boldsymbol{\xi}(t)$$

$$\frac{1}{2} \langle \boldsymbol{\xi}(t) \boldsymbol{\xi}(s) + \boldsymbol{\xi}(s) \boldsymbol{\xi}(t) \rangle_{\text{bath}} = \\ \cdot \frac{m}{\pi} \int_0^\infty \text{Re} \hat{\gamma}(-i\omega + 0^+) \hbar\omega \coth \left(\frac{\hbar\omega}{2k_B T} \right) \cos [\omega(t-s)] d\omega$$

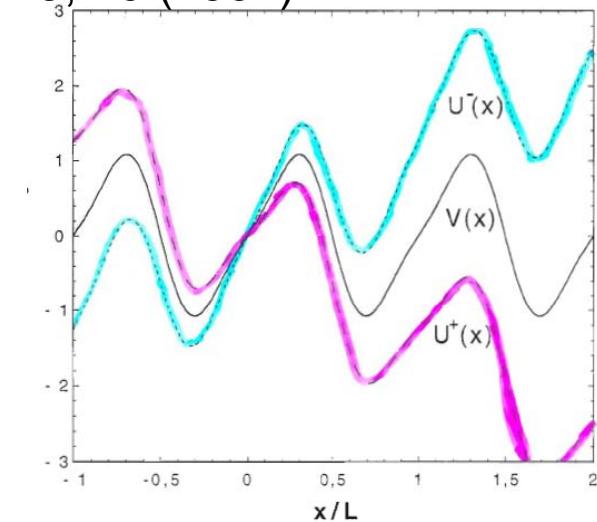
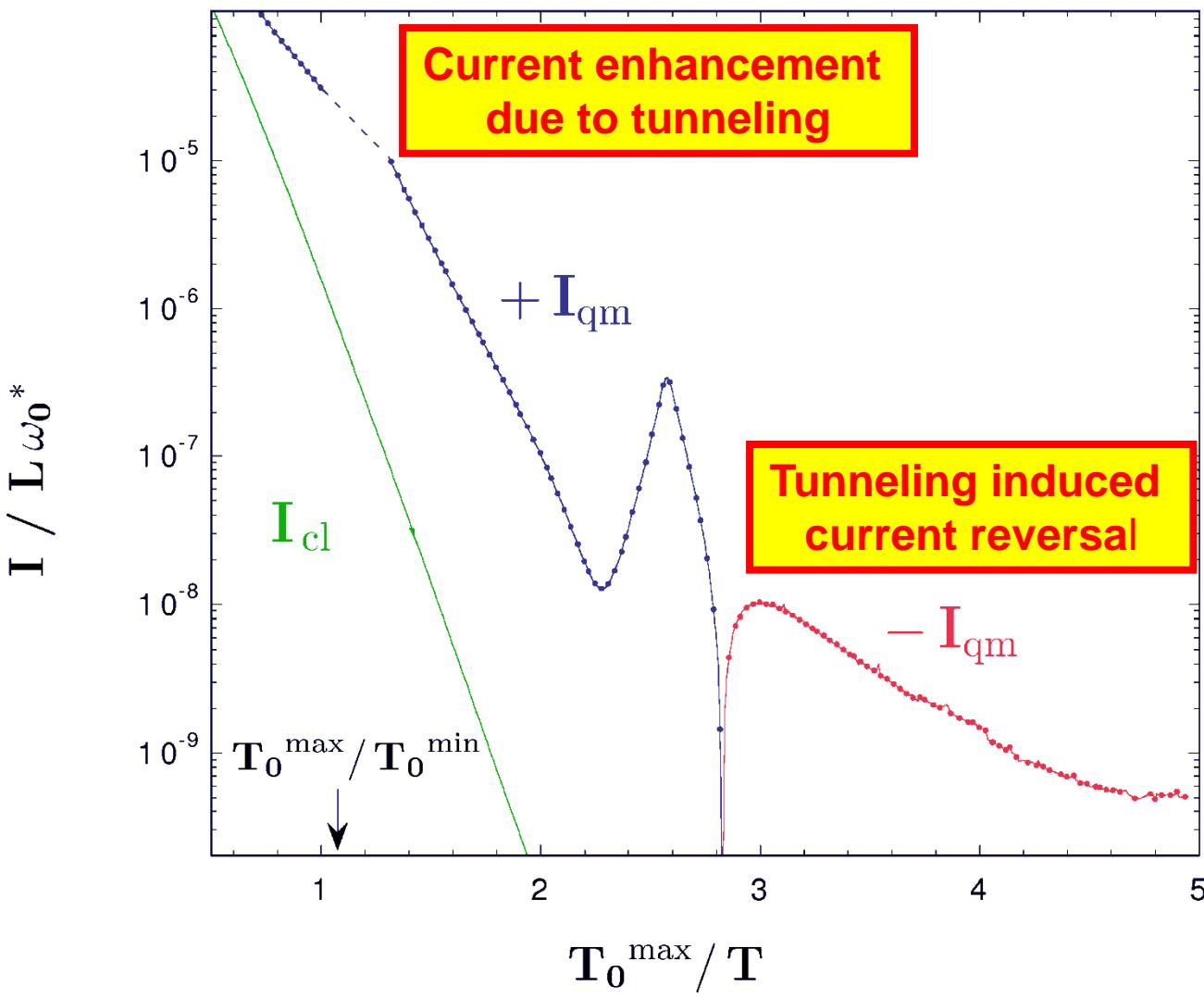
And:

$$[\boldsymbol{\xi}(t), \boldsymbol{\xi}(s)] = -i\hbar \dots \neq 0$$



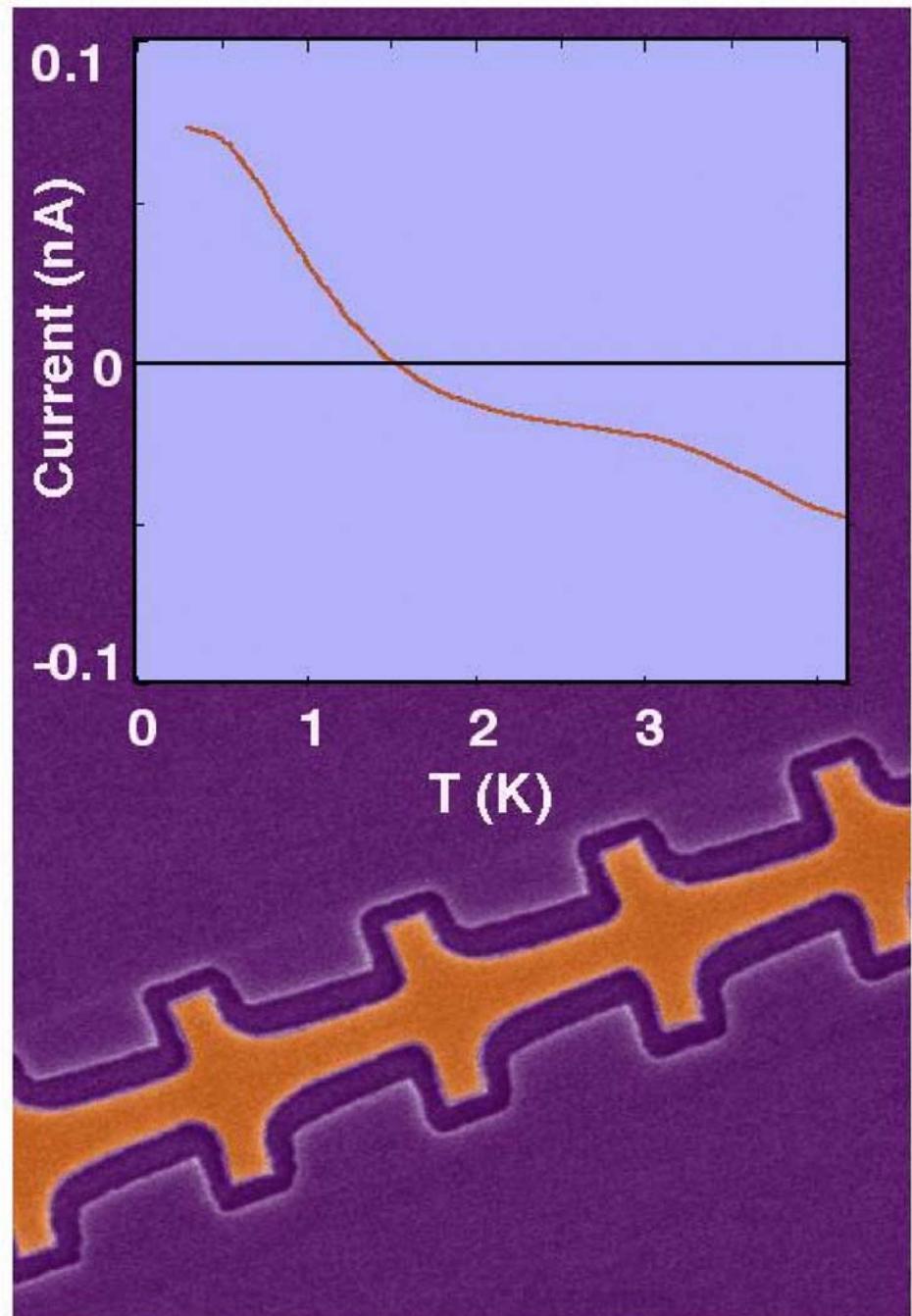
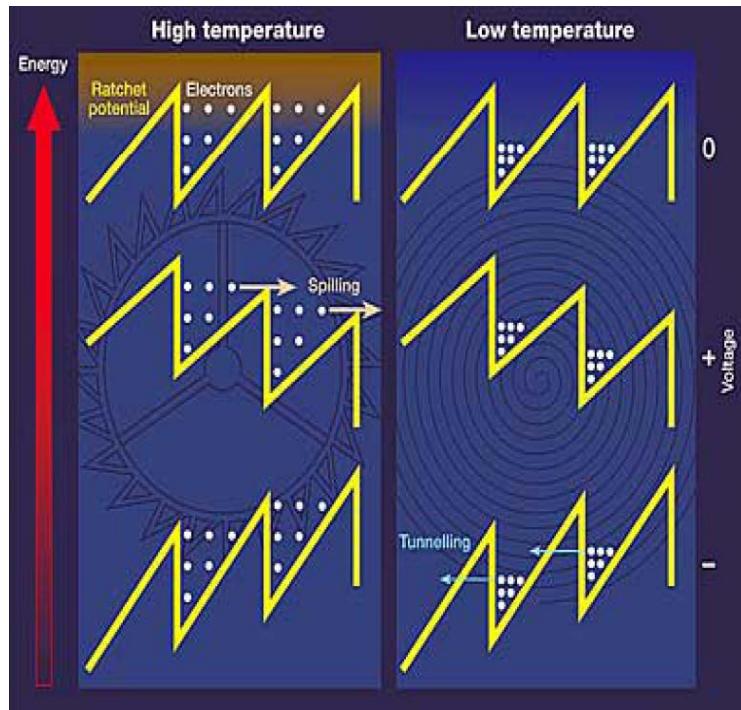
Rocking Ratchet - Theory

P. Reimann, M. Grifoni, P. H., Phys. Rev. Lett. **79**, 10 (1997)

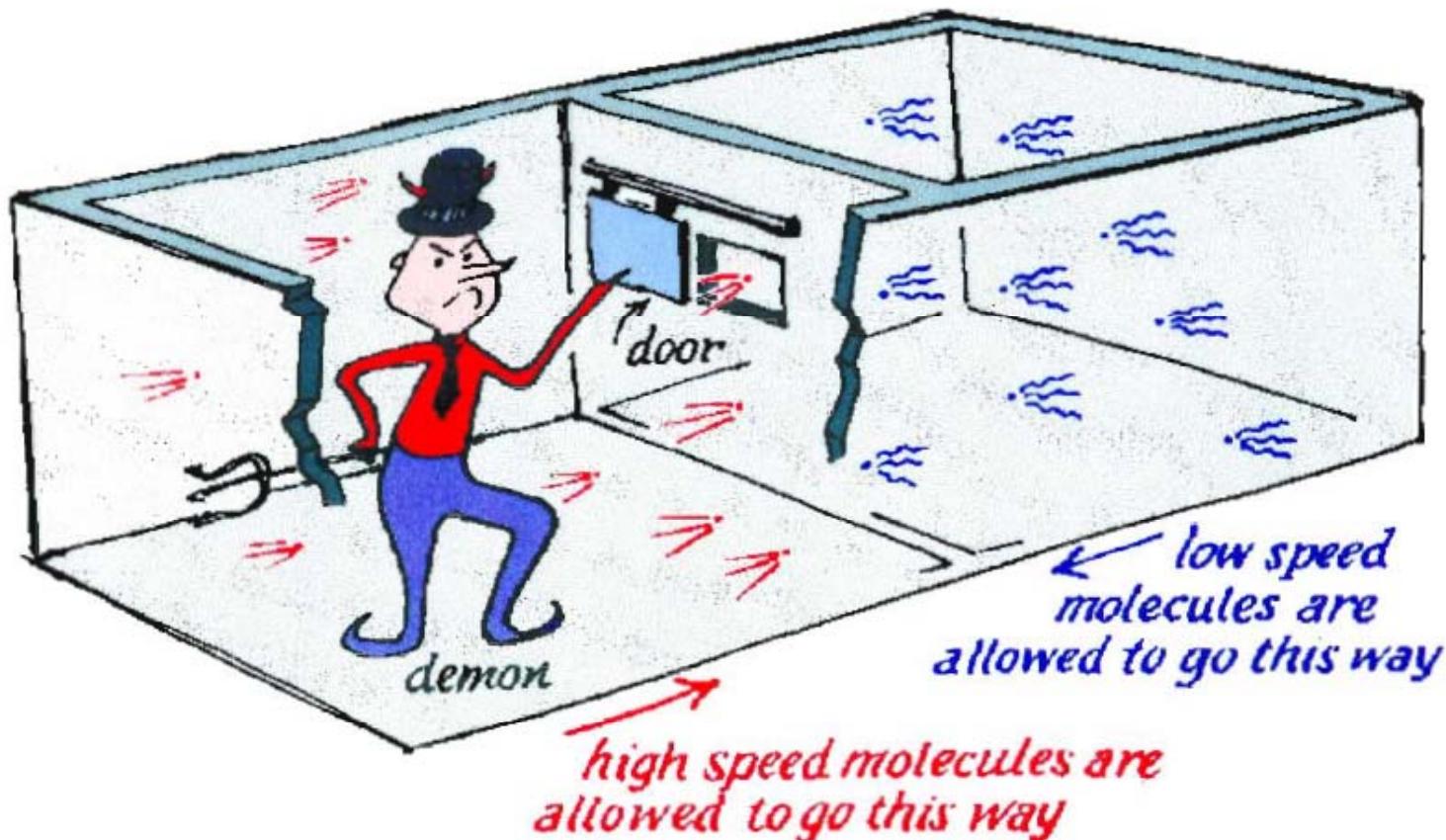


Rocking QM Ratchet – Experiment

H. Linke, *et al.*,
SCIENCE 286, 2314 (1999)



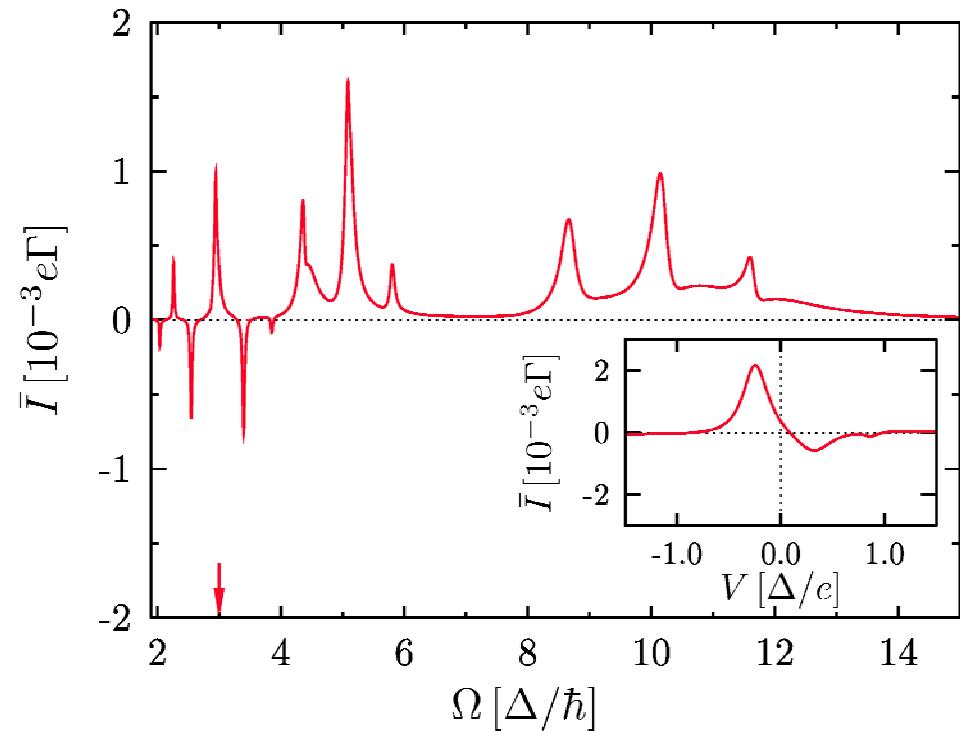
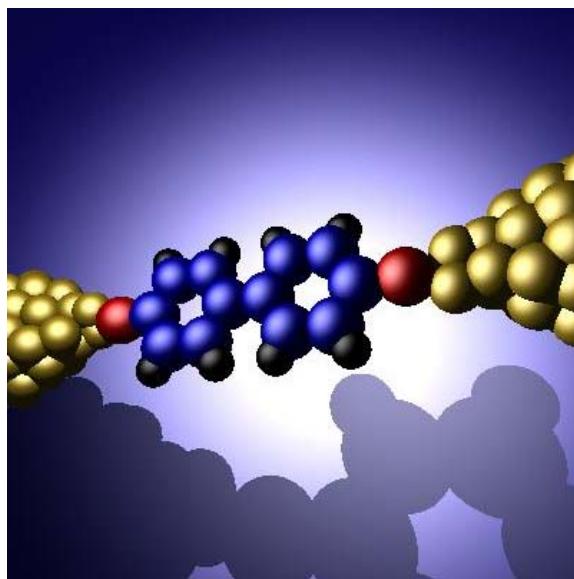
Brownian motors: EX(E/O)RCISING DEMONS



Source: H.S. Leff, *Maxwell's Demon* (Adam Hilger, Bristol, 1990)

Molecular wires

J. Lehmann, S. Kohler, P. H., A. Nitzan, Phys. Rev. Lett. **88**, 228305 (2002)

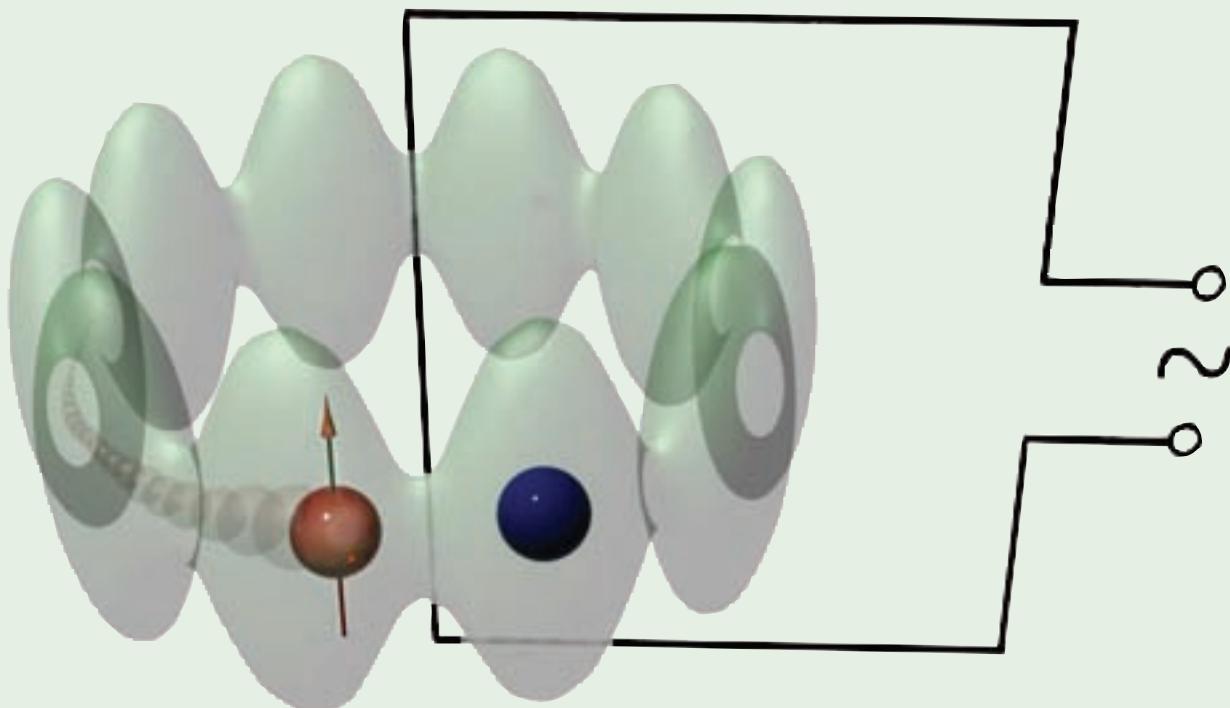


- Field strength $E=106$ V/cm
- $\Omega=3\Delta$ corresponds to $4\mu\text{m}$ wavelength
- typical current: some nA

PHYSICAL REVIEW LETTERS

Member Subscription Copy
Library or Other Institutional Use Prohibited Until 2014

Articles published week ending 12 JUNE 2009



Published by the
American Physical Society

APS
physics

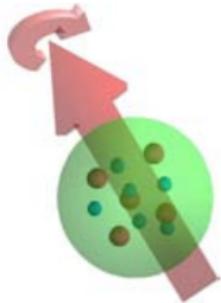
Volume 102, Number 23

Quantum Gears: Driving Through Interactions

Area:

Quantum-Nanophysics;
some scientific
visions

Peter Hänggi



cold atoms (bosons, fermions)

$$\tilde{H} = H_R + H_B(t) + H_{RB}, \quad H_B(t+T) = H_B(t)$$

- targeting excitations
- experimentally feasible setup
- full many-body treatment

Generalizations of Brownian Motion

Brownian motion: Generalized Langevin-equation

Hamiltonian: $H_{\text{tot}} = H_{\text{sys}} + H_{\text{env}} + H_{\text{ww}}$

→ $m\ddot{\mathbf{x}}(t) + \int_{-\infty}^t \gamma(t-t') \dot{\mathbf{x}}(t') dt' + V'(\mathbf{x}; t) = \xi(t)$

Asymptotically normal, anomalously fast, or anomalously slow
– via fractional Brownian motion –

$$\int_0^\infty \gamma(t) dt = \begin{cases} \text{const} & \Rightarrow \text{normal} \\ 0 & \Rightarrow \text{superfast} \\ \infty & \Rightarrow \text{superslow} \end{cases}$$

Connection to the fractional Fokker-Planck-equation

Fractional Fokker-Planck equation

Subdiffusion ($\alpha < 1$):

$$\frac{\partial P(x,t)}{\partial t} = \left[\frac{\partial}{\partial x} \frac{V'(x,t)}{\gamma_\alpha} + K_\alpha \frac{\partial^2}{\partial x^2} \right] {}_0D_t^{1-\alpha} P(x,t)$$

Fat tails in the distribution of the residence times

Riemann-Liouville Operator



Superdiffusion ($\alpha > 1$):

$$\frac{\partial P(x,t)}{\partial t} = \left[\frac{\partial}{\partial x} \frac{V'(x,t)}{\gamma_\alpha} + K_\alpha \frac{\partial^{2/\alpha}}{\partial |x|^{2/\alpha}} \right] P(x,t)$$

Fat tails in the distribution of the jump lengths

Riesz-derivative

Fractional Fokker-Planck equation

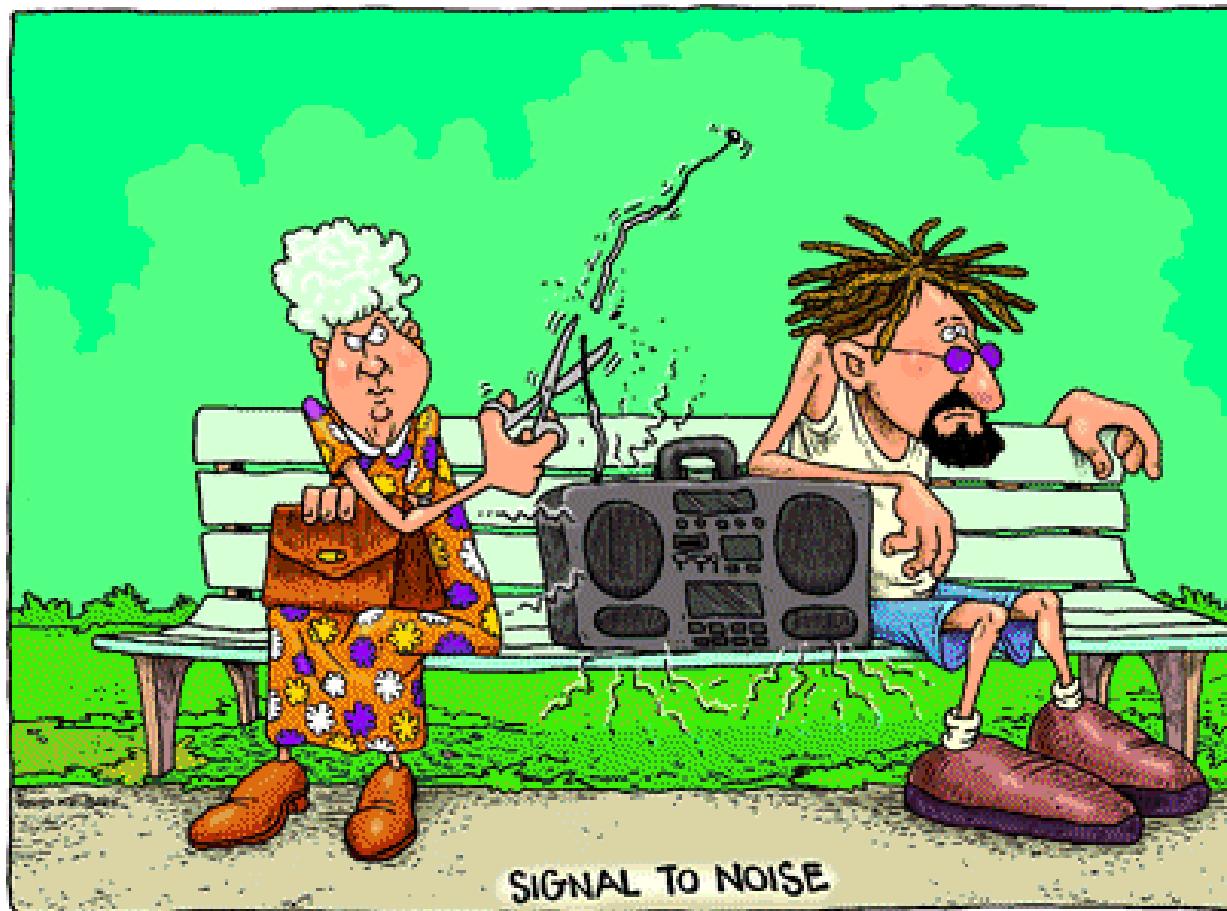
subdiffusive ($\alpha < 1$)

$$\frac{\partial P(x, t)}{\partial t} = \left[\frac{\partial}{\partial x} \frac{V'(x, t)}{\eta_\alpha} + K_\alpha \frac{\partial^2}{\partial x^2} \right] {}_0 D_t^{1-\alpha} P(x, t)$$

Riemann-Liouville Operator

$${}_0 D_t^{1-\alpha} f(t) = \frac{1}{\Gamma(\alpha)} \frac{d}{dt} \int_0^t \frac{f(t')}{(t-t')^{1-\alpha}} dt'$$

Noise – always bad ?



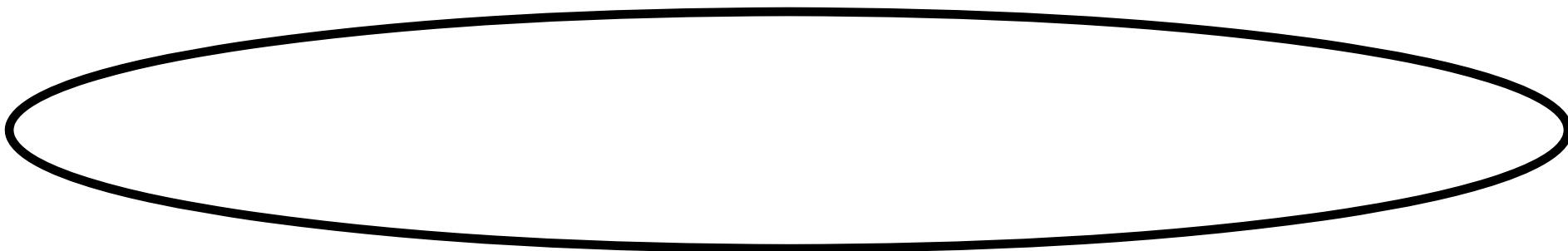
Source: Agilent Technologies

Brownian motion

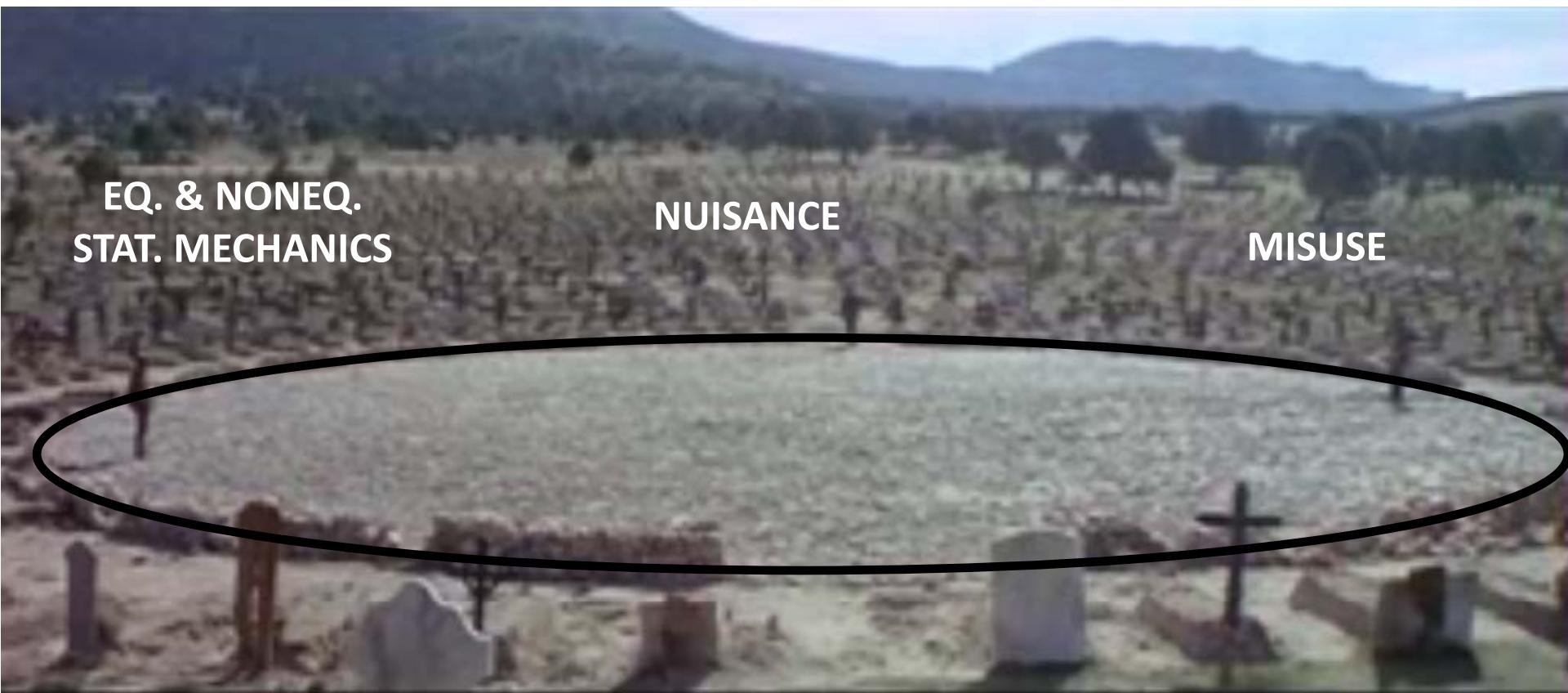
EQ. & NONEQ.
STAT. MECHANICS

NUISANCE

MISUSE



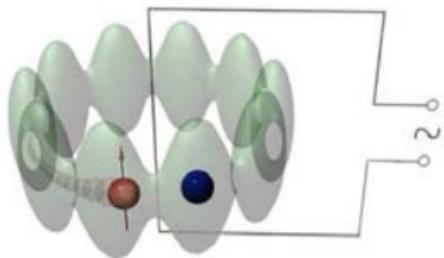
The good, the bad and the simply silly



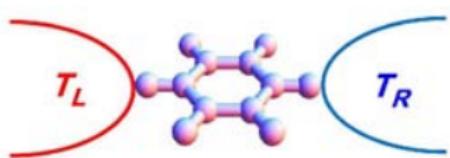
Driven Many-Body Quantum Systems

Area:
Quantum-
Nanophysics;
some scientific
visions

Peter Hänggi



- quantum machinery



- role of heat transport in nanodevices

What “is” thermodynamics ?

- ✓ non-local description in terms of **symmetry** (breaking) parameters

“Good” starting point in relativistic thermodynamics ?

- ✓ (non-)conserved tensor densities, **Noether currents**

Origin of different temperature transformation laws ?

- ✓ choice of space-time hyperplanes
- ✓ definition of heat, formulation of 1st/2nd law

How should one define thermodynamic **observables** in special and general relativity?

- ✓ invariant manifolds, **lightcone integrals**

Observable consequence:
temperature-induced apparent drift

