## Comment on "Experimental Verification of a Jarzynski-Related Information-Theoretic Equality by a Single Trapped Ion"

Reference [1] reports on the experimental verification of an identity in probability theory that reads [see Eq. (3) of Ref. [1]]:

$$\langle e^{-I_{nm}} \rangle = \sum_{mn} p_{nm} e^{-I_{nm}} = 1.$$
 (1)

It is claimed in Ref. [1] that Eq. (1) implies Eq. (4) below via the relation [see Eq. (4) of Ref. [1]]:

$$I_{nm} = \beta (E'_m - E_n - F' + F) \quad (\text{not correct}).$$
(2)

Our comment is that while (under specific conditions) Eqs. (1), (4) can be simultaneously valid, it is not true that the quantities appearing in the exponents are the same. Equation (2) is not valid. Accordingly, it is not correct to state that Eq. (1) is related to, or implies Eq. (4) [2].

Equation (1) presents an identity that holds for any joint probability  $p_{nm}$ . It follows from the Bayes rule  $p_{nm} = p_{m|n}p_n$ , normalization  $\sum_{nm}p_{nm} = \sum_n p_n = \sum_m q_m = 1$  and the definition  $I_{nm} = \ln p_{m|n} - \ln q_m$  in [1]:

$$\sum_{mn} p_{nm} e^{-I_{nm}} = \sum_{mn} \frac{p_{nm} q_m}{p_{m|n}} = \sum_n p_n \sum_m q_m = 1.$$
(3)

Equation (2) is claimed to be valid for all quantum systems evolving according to a CPTP map, starting in thermal equilibrium and being subject to a two-point energy measurement, for which  $p_{nm} = \text{Tr}Q_m \sum_i \Lambda_i (P_n \rho P_n) \Lambda_i^{\dagger} Q_m$  [1], "if the system is initially prepared as a Gibbs state" [1]. We first point out that if that were true, upon inserting Eq. (2) in Eq. (1), it would imply that

$$\langle e^{-\beta(E'_m - E_n - F' + F)} \rangle = 1. \tag{4}$$

This contradicts the well-known fact that, under those conditions, rather the following holds true [3-8]:

$$\langle e^{-\beta(E'_m - E_n - F' + F)} \rangle = \gamma = \sum_i \operatorname{Tr} \Lambda_i^{\dagger} \rho \Lambda_i$$
 (5)

with  $\gamma$  being generally different from 1 [9].

That Eq. (2) is not valid can be checked by direct inspection. For  $I_{nm}$  we find

$$I_{nm} = \ln \frac{p_{m|n}}{q_m} = \ln \frac{p_{nm}}{p_n q_m} = \ln \frac{\text{Tr}Q_m \sum_i \Lambda_i P_n \rho P_n \Lambda_i^{\dagger}}{(\text{Tr}P_n \rho)(\text{Tr}Q_m \sum_i \Lambda_i \rho \Lambda_i^{\dagger})}$$
$$= \ln \frac{(\sum_k e^{-\beta E_k})(\text{Tr}Q_m \sum_i \Lambda_i P_n \Lambda_i^{\dagger})}{\text{Tr}Q_m \sum_{i,k} e^{-\beta E_k} \Lambda_i P_k \Lambda_i^{\dagger}}$$
(6)

where we have used  $\sum_{i} \Lambda_{i}^{\dagger} \Lambda_{i} = 1$ ,  $\rho = e^{-\beta H}/Z = \sum_{k} P_{k} e^{-\beta E_{k}}/Z$ ,  $P_{k} P_{n} = \delta_{kn} P_{n}$ ,  $q_{m} = \sum_{n} p_{nm} = \operatorname{Tr} Q_{m} \sum_{i,n} \Lambda_{i} P_{n} \Lambda_{i}^{\dagger} e^{-\beta E_{n}}/Z$ , and  $p_{n} = \sum_{m} p_{nm} = \operatorname{Tr} P_{n} \rho = e^{-\beta E_{n}}/Z$  [1]. For  $\beta(E'_{m} - E_{n} - F' + F)$  we find

$$\beta(E'_m - E_n - F' + F) = \ln \frac{(\sum_k e^{-\beta E'_k}) e^{\beta(E'_m - E_n)}}{\sum_k e^{-\beta E_k}}.$$
 (7)

Note that the final eigenvalues  $E'_k$  enter explicitly in Eq. (7) while Eq. (6) is *independent* of the final eigenvalues  $E'_k$ . Similarly, the projectors  $P_k$ ,  $Q_m$  enter Eq. (6) explicitly whereas they do not appear in Eq. (7). Therefore the two quantities are never equal. This argument remains valid as well for the special case of an unitary evolution,  $U_C$ , that commutes with the  $P_k$ 's, as employed in Ref. [1], for which we obtain

$$I_{nm} = \ln \frac{(\sum_{k} e^{-\beta E_{k}})(\operatorname{Tr} Q_{m} P_{n})}{\operatorname{Tr} Q_{m} \sum_{k} e^{-\beta E_{k}} P_{k}}.$$
(8)

In that case Eq. (1) and Eq. (4) hold simultaneously (as verified experimentally in Ref. [1]) but Eq. (2) is, however, nevertheless *not* valid.

Michele Campisi<sup>1</sup> and Peter Hänggi<sup>2</sup> <sup>1</sup>Dipartimento di Fisica e Astronomia Università di Firenze and INFN Sezione di Firenze Via G. Sansone 1, I-50019 Sesto Fiorentino (FI), Italy <sup>2</sup>Institute of Physics University of Augsburg Universitätsstraße 1, D-86135 Augsburg, Germany

Received 29 January 2018; published 24 August 2018 DOI: 10.1103/PhysRevLett.121.088901

- [1] T. P. Xiong, L. L. Yan, F. Zhou, K. Rehan, D. F. Liang, L. Chen, W. L. Yang, Z. H. Ma, M. Feng, and V. Vedral, Phys. Rev. Lett. **120**, 010601 (2018).
- [2] We are not using the notation  $\Delta F = F F'$  employed in Ref. [1], which has opposite sign convention as compared to the standard notation [10], in order to minimize confusion. We also refrain from using the notation  $W = E_n E'_m$  of [1] which would hint at the fact that the energy change in the system can be interpreted as work. This is in fact generally not the case, in particular, it is not the case for open quantum systems (a situation that can in certain cases be described by CPTP maps), in which case the change in a system energy generally contains both work and heat [10,11].
- [3] Y. Morikuni and H. Tasaki, J. Stat. Phys. 143, 1 (2011).
- [4] D. Kafri and S. Deffner, Phys. Rev. A 86, 044302 (2012).
- [5] A. E. Rastegin, J. Stat. Mech. (2013) P06016.
- [6] T. Albash, D. A. Lidar, M. Marvian, and P. Zanardi, Phys. Rev. E 88, 032146 (2013).
- [7] G. Watanabe, B. P. Venkatesh, P. Talkner, M. Campisi, and P. Hänggi, Phys. Rev. E 89, 032114 (2014).

- [8] M. Campisi, J. Pekola, and R. Fazio, New J. Phys. 19, 053027 (2017).
- [9] A condition for  $\gamma$  to equal 1 is that the dynamics be unitary, or more generally, unital [3–8]; i.e.,  $\sum_i \Lambda_i \Lambda_i^{\dagger} = 1$ .
- [10] M. Campisi, P. Hänggi, and P. Talkner, Rev. Mod. Phys. 83, 771 (2011); 83, 1653 (2011).
- [11] P. Talkner, M. Campisi, and P. Hänggi, J. Stat. Mech. (2009) P02025.